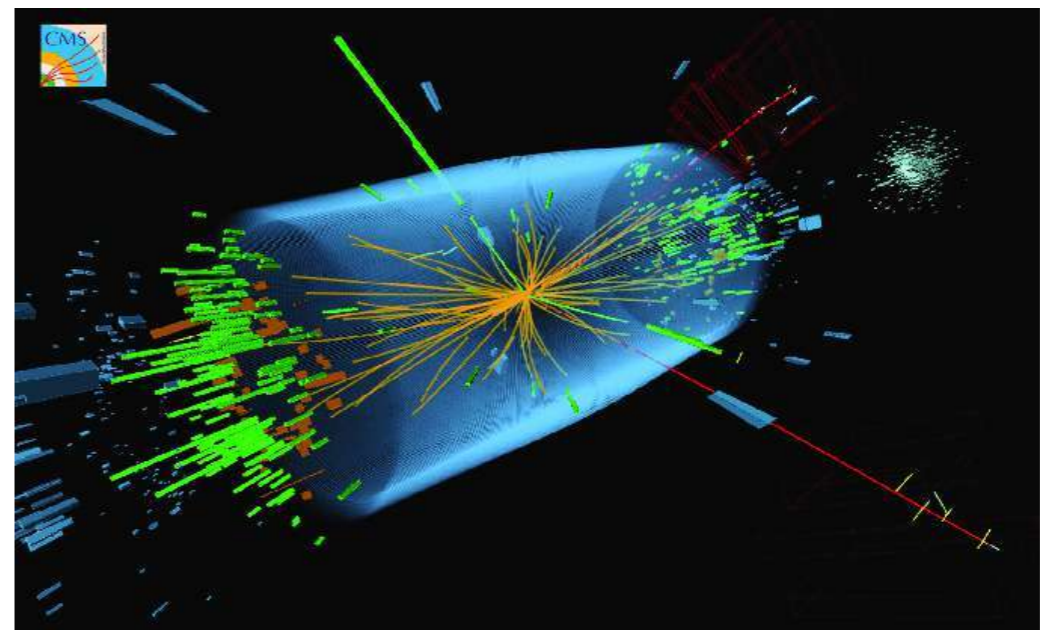
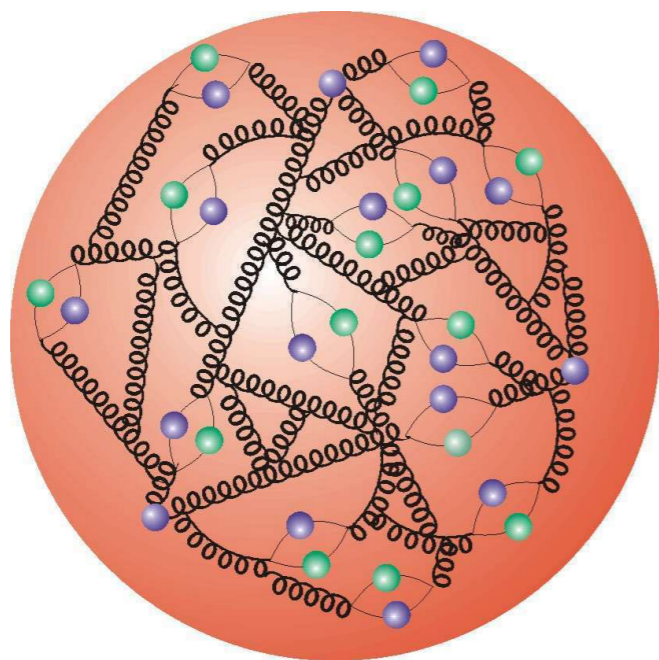


From Protons to Collisions...

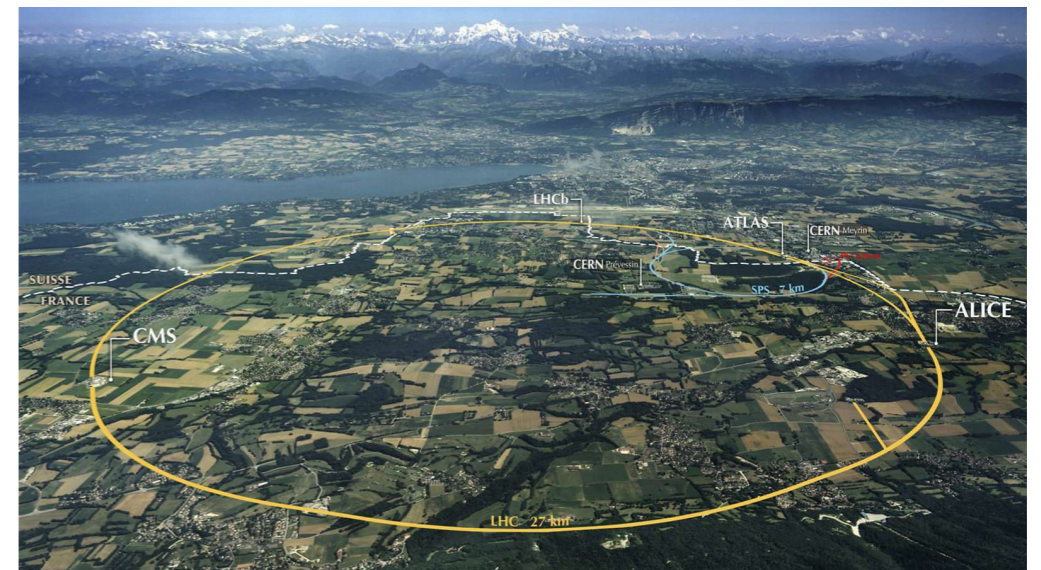
Lucian Harland-Lang, University of Oxford

Saturday Mornings of Theoretical Physics, Oxford,
May 10 2019



The LHC: a proton-proton collider

- The **LHC** works by colliding proton beams head on at high energy.
- We examine the debris of these interactions for signs of the Higgs and its interactions.
- Before getting to that: we need to understand what we are colliding.
- The **proton**: What is it? What is it made of? How does it behave when given the LHC treatment?

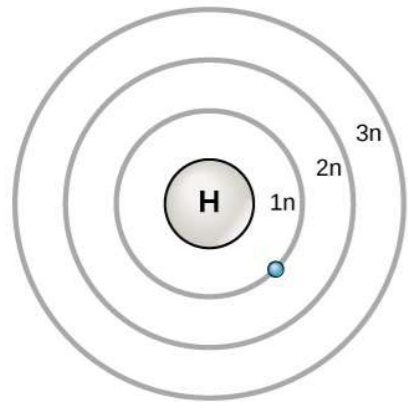


What is the Proton?

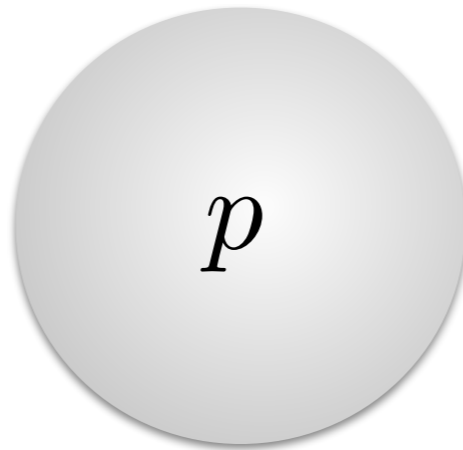
A Proton Roadmap

- In 20th century, layers of proton complexity/substructure uncovered:

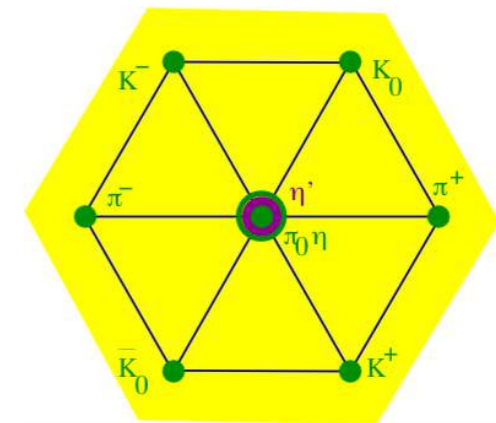
1910s: Rutherford and the compact atomic nucleus



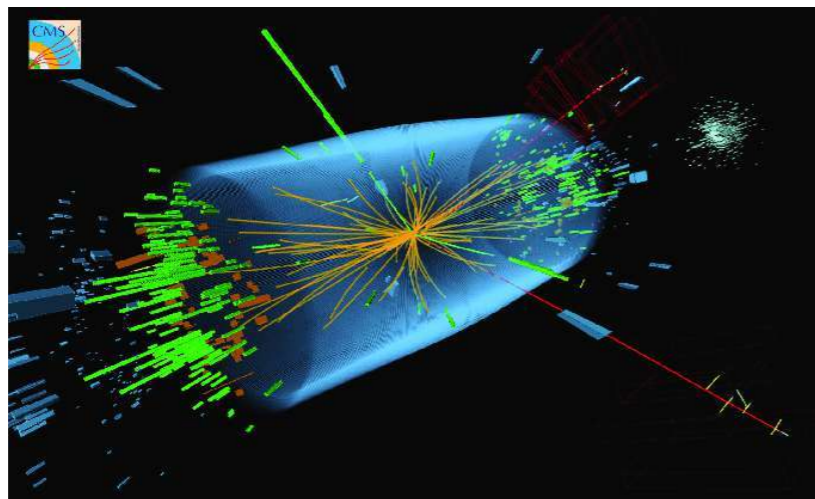
1950s: Measurement of proton radius



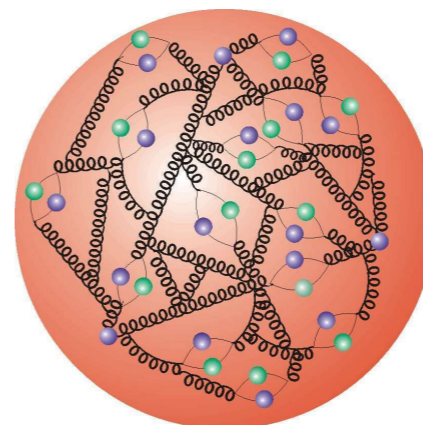
1960s: The Quark Model



2019: proton collisions, the LHC discovery tool!

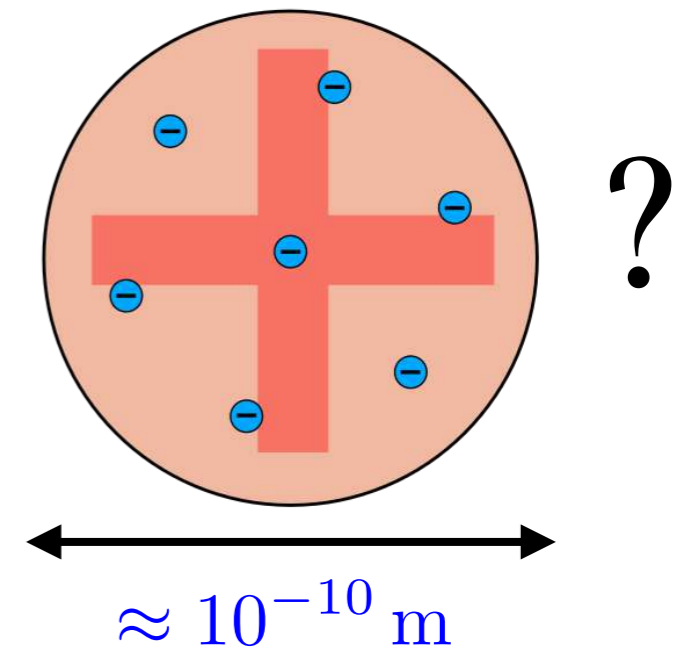


1970s: QCD and the proton.



Observing Nuclear Structure

- Start of 20th century: no clear picture of structure of atom
- Need to 'look' **inside atom**. Idea: fire beam of particles at target of interest, and measure the corresponds scattering in detector.

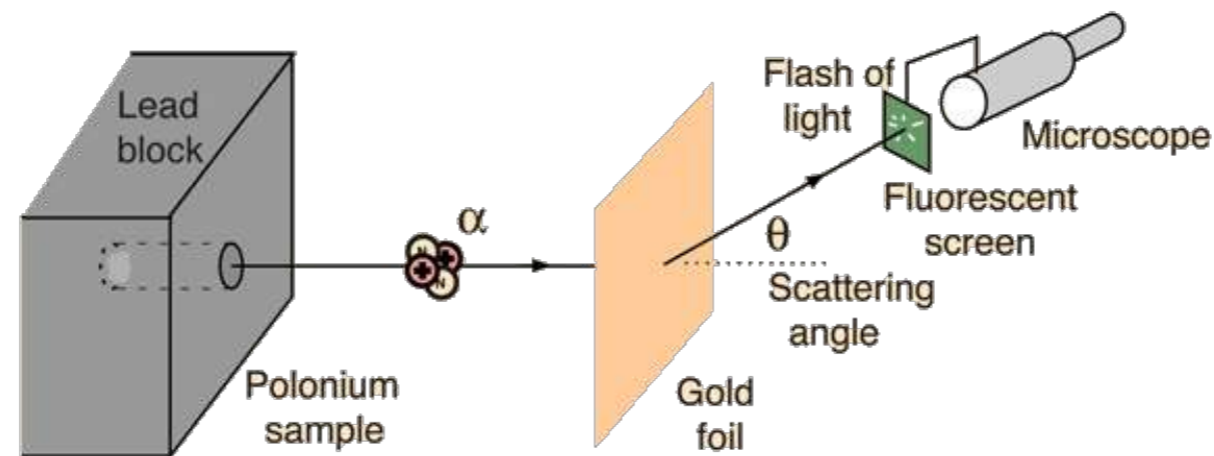


Beam: alpha particles from radioactive source

Target: Gold Foil

Detector: Geiger and Marsden sitting in basement

**Rutherford
in 1909:**



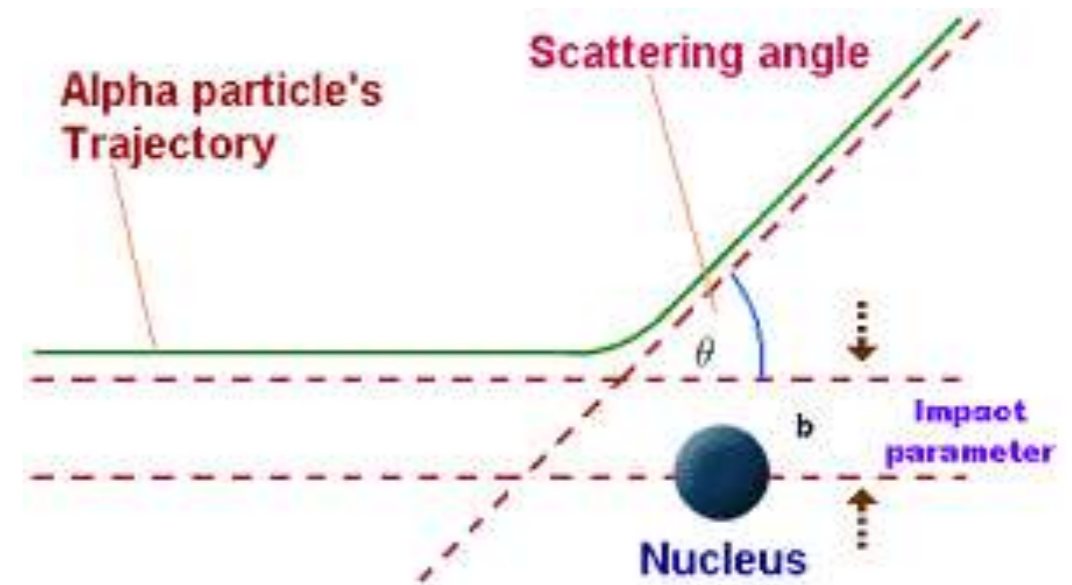
- This basic idea is fundamentally unchanged today: forms the **basis of all particle physics**.

Rutherford Scattering

- α particles observed to occasionally deflect at **very large angles**. Could only be explained by compact, \sim point-like nucleus!

- Data well described by scattering due to electrostatic repulsion by positive **'point-like'** charge:

$$\frac{dN}{d \cos \theta} \propto \frac{Z^2 \alpha^2}{\sin^4 \left(\frac{\theta}{2} \right)}$$



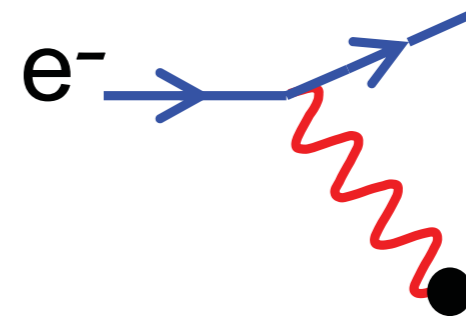
- In what sense was nucleus point-like? Only within **resolution of experiment**.
- Here 'point-like' scattering observed up to closest permitted approach of alpha particles **~ 30 fm** ($\sim 1/10000$ th of atomic radius).

$$1 \text{ fm} = 10^{-15} \text{ m}$$

Looking Into the Proton

- By 1950s still no idea of internal structure of proton/neutrons. How to probe this? Use cleaner electron beam.

- Electron interacts with proton via photon exchange. **Resolution** \leftrightarrow **photon wavelength**:



$$E = hf \Rightarrow \lambda = \frac{hc}{E}$$

- Thus **higher energies** \Rightarrow see **further inside** proton! Plugging in some (rough!) numbers...

\sim Proton size

Energy unit (LHC beam = 7000 GeV)

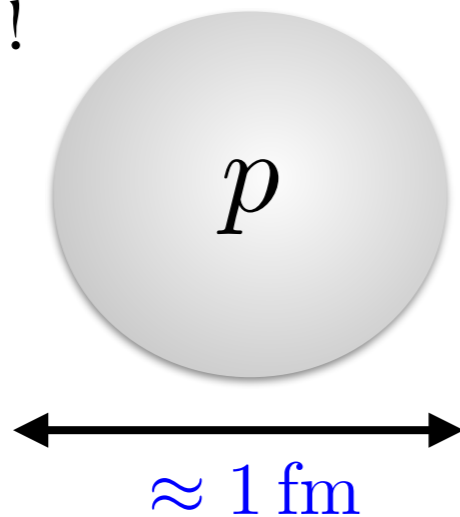
$$\lambda \sim 1 \text{ fm} \Rightarrow E = 1 \text{ GeV}$$

- Much less than LHC energies but at time required cutting edge technology to accelerate electrons.

- 1950s experiments of **Hofstadter**: collisions of electron beam (up to ~ 0.5 GeV) with protons.

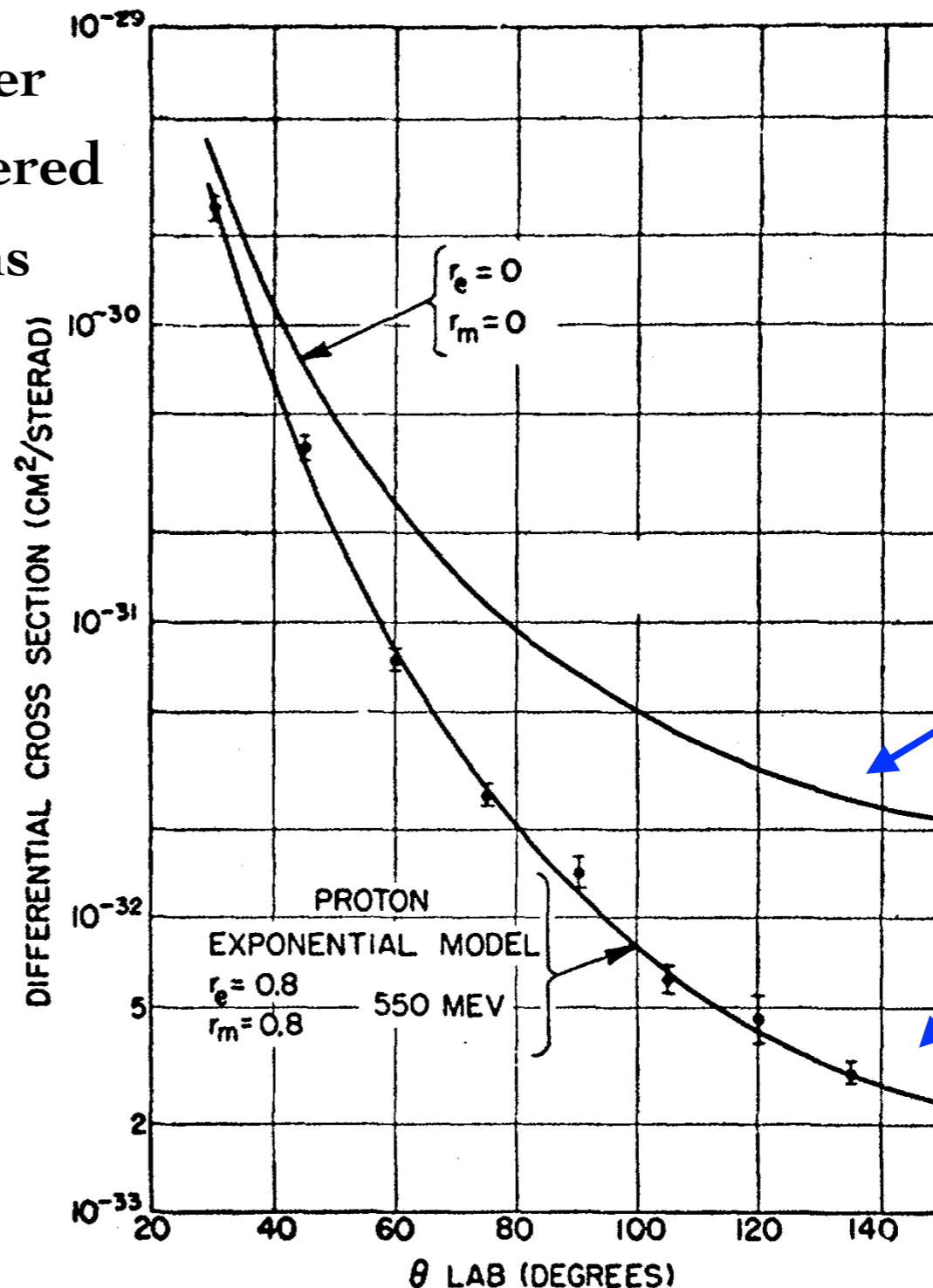


- **Clear deviation** from point-like behaviour seen \Rightarrow **proton has size!**



\sim Number of scattered electrons

- What is the nature of this extended object? Continuous extent of charged 'stuff', or formed of something **more fundamental?**



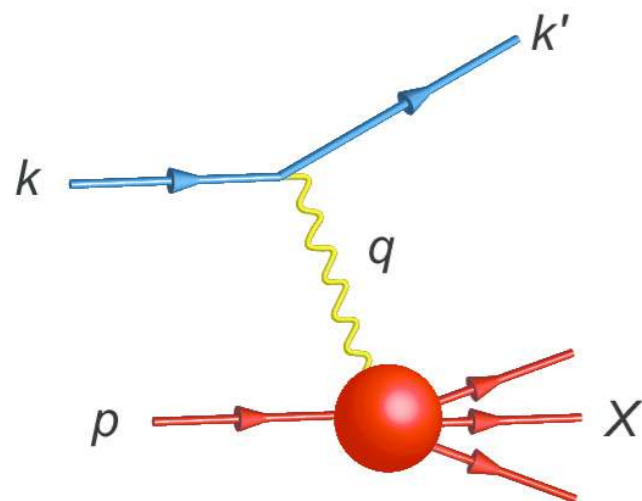
$$\frac{dN}{d \cos \theta} \propto \frac{Z^2 \alpha^2}{\sin^4 \left(\frac{\theta}{2} \right)}$$

Point-like theory

Data

Looking Deeper

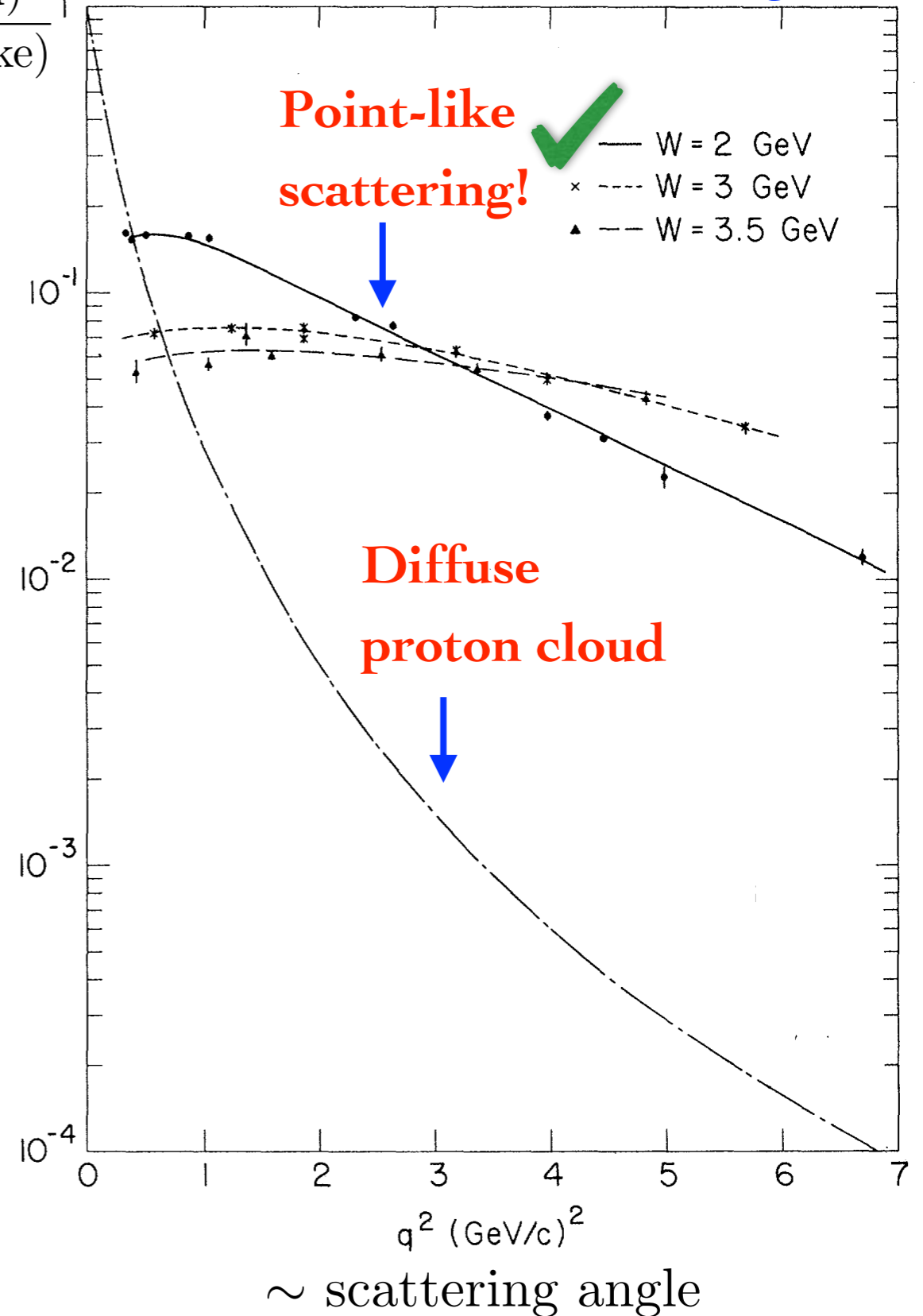
- **1976**: The Stanford Linear Electron Collider (**SLAC**) collider pushed to unprecedented electron energies.
- Electrons scattered inelastically, destroying proton, seeing deeper inside.
- Something **completely new**: scattering of electron with **point-like object**!



‘Deep Inelastic Scattering’

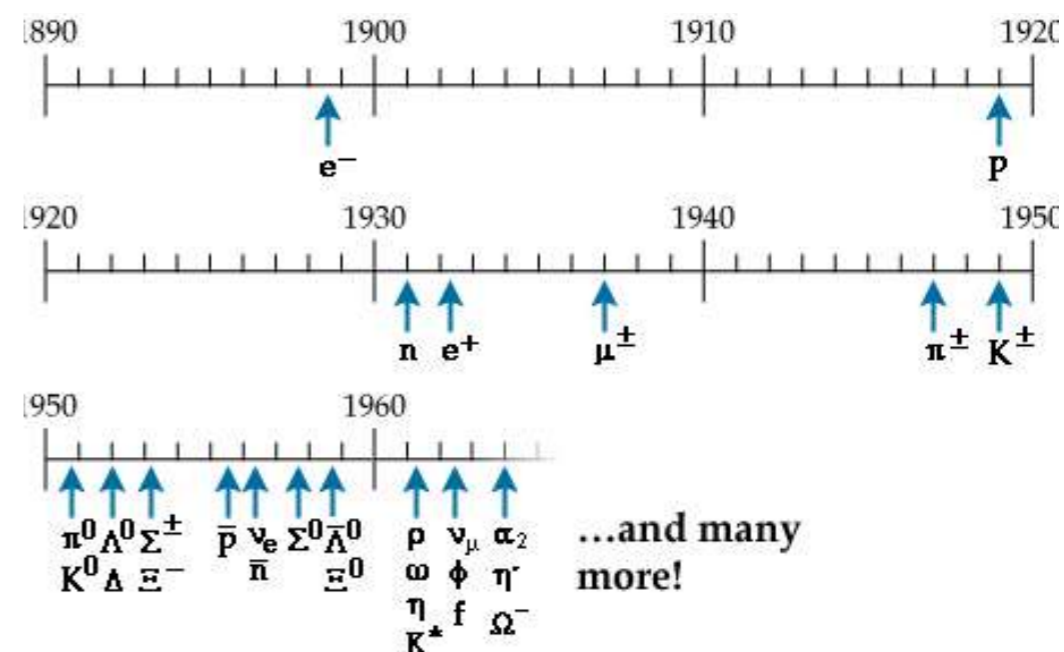
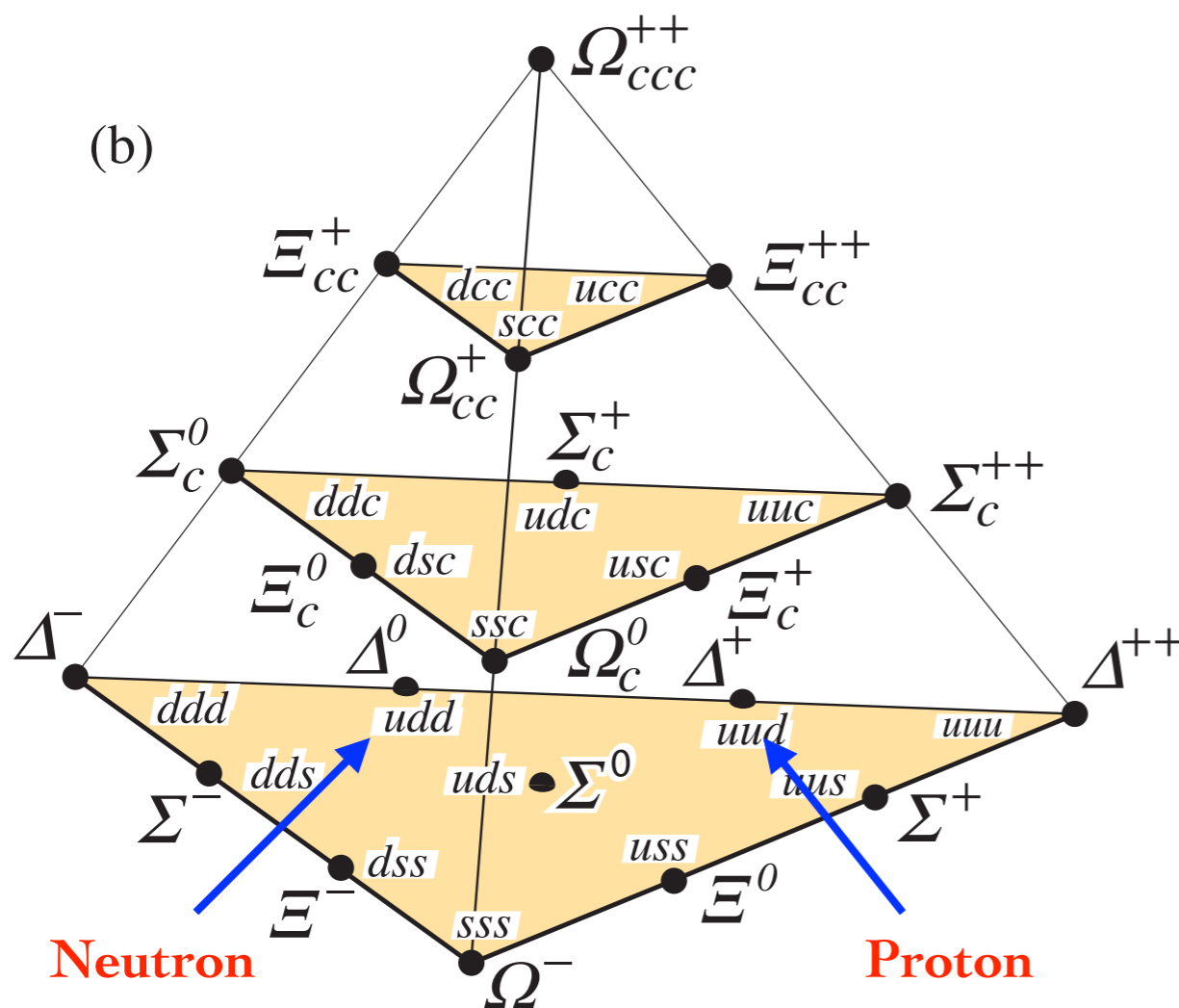
$$\frac{\sigma(\text{observed})}{\sigma(\text{point - like})}$$

‘Cross Section’ $\sigma \sim$ Scattering Rate



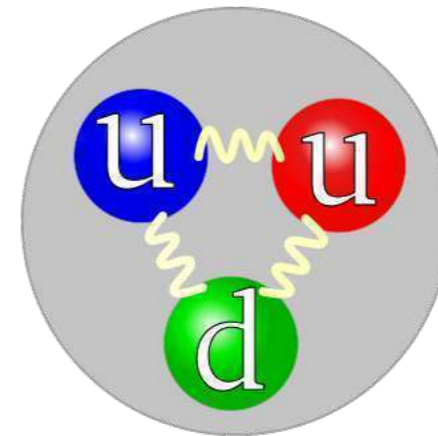
Quarks

- These objects, first given the agnostic label of ‘**partons**’ by Feynman, but soon identified with the ‘**quarks**’ needed to explain the ‘**hadron zoo**’.
- Initially required 3 distinct types (**up, down, strange**). Since then 3 more discovered (**charm, bottom, top**).



The Proton

- **Proton** naturally described in quark model: one part of a larger family of 'baryons' - **bound state of 3 quarks**.
- Proton: two 'up' quarks and one 'down'.
- Quarks carry fractional electric charge:



Up quark: $+2/3$

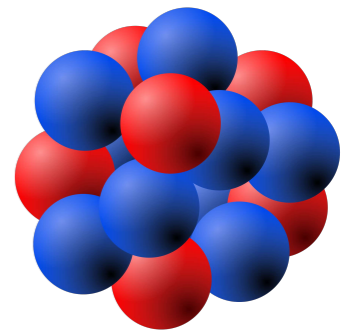
Down quark: $-1/3$

Proton: $2 \times \frac{2}{3} - \frac{1}{3} = +1$ ✓

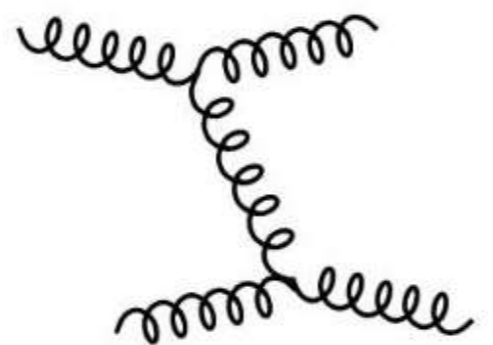
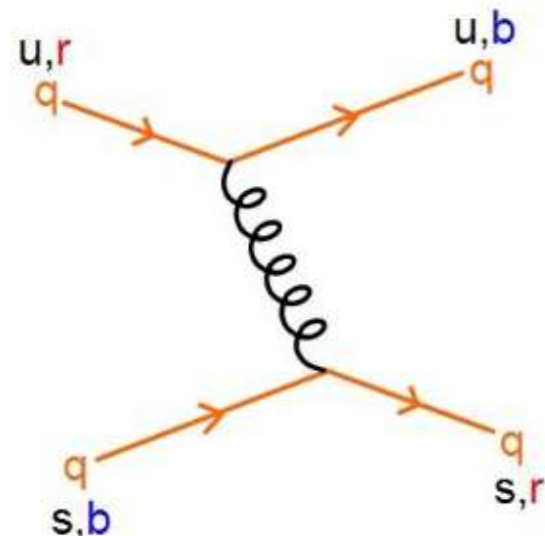
- This is the basic idea, and it still holds true today. But more fundamental questions remained unanswered:

- ❖ What is the **force** that is binding these quarks together?
- ❖ How do we describe it theoretically?

Quantum Chromodynamics



- Nuclei as bound states of proton and neutrons: required introduction of completely new attractive **binding force**: the '**strong nuclear force**'.
- Can now be understood in terms of more fundamental quark interactions.
- New force carrying '**gluon**' particle mediating interactions between quarks, which carry '**colour**' charge.
- Mathematical framework direct generalisation of QED. New effects:
 - ❖ **3** different colour charges (just 1 electric charge).
 - ❖ Gluon carries charge: can **self-interact!**

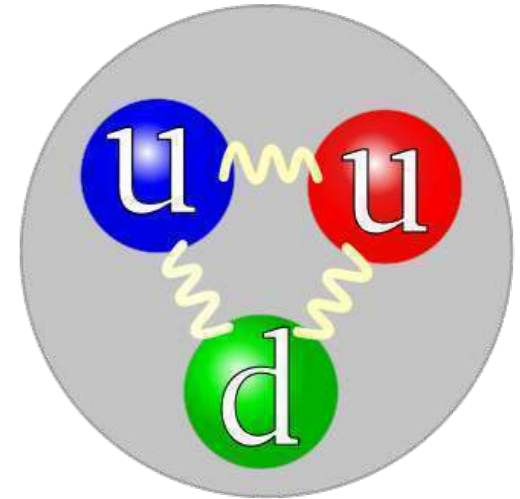


gluon-gluon scattering

QED	↔	QCD
Electrons	↔	Quarks
Photons	↔	Gluons
Electric Charge	↔	Colour Charge

**What do proton collisions at the
LHC look like?**

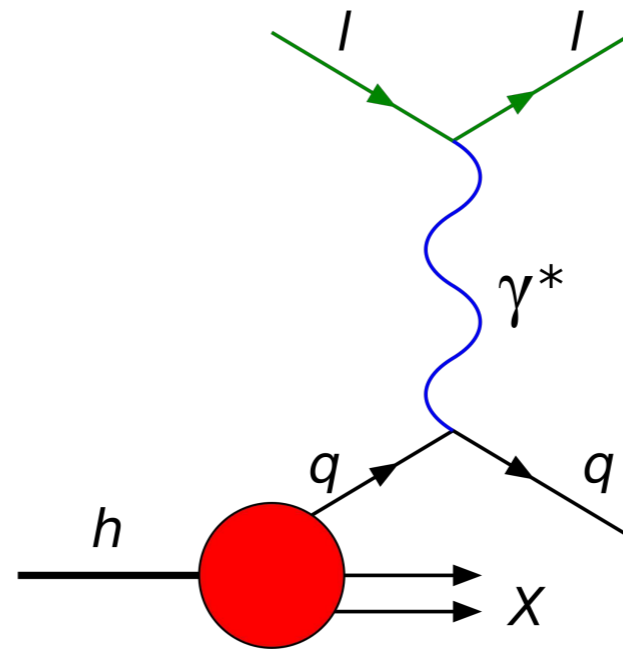
Colliding Protons



- How do we apply this model to proton collisions at the **LHC**?
- For now, stick with simpler electron-proton collision case.
- Basic idea: potentially complex **electron-proton** interaction really due to simple scattering between point-like **electron** and **quark** within proton.

**Scattering
Rate:**

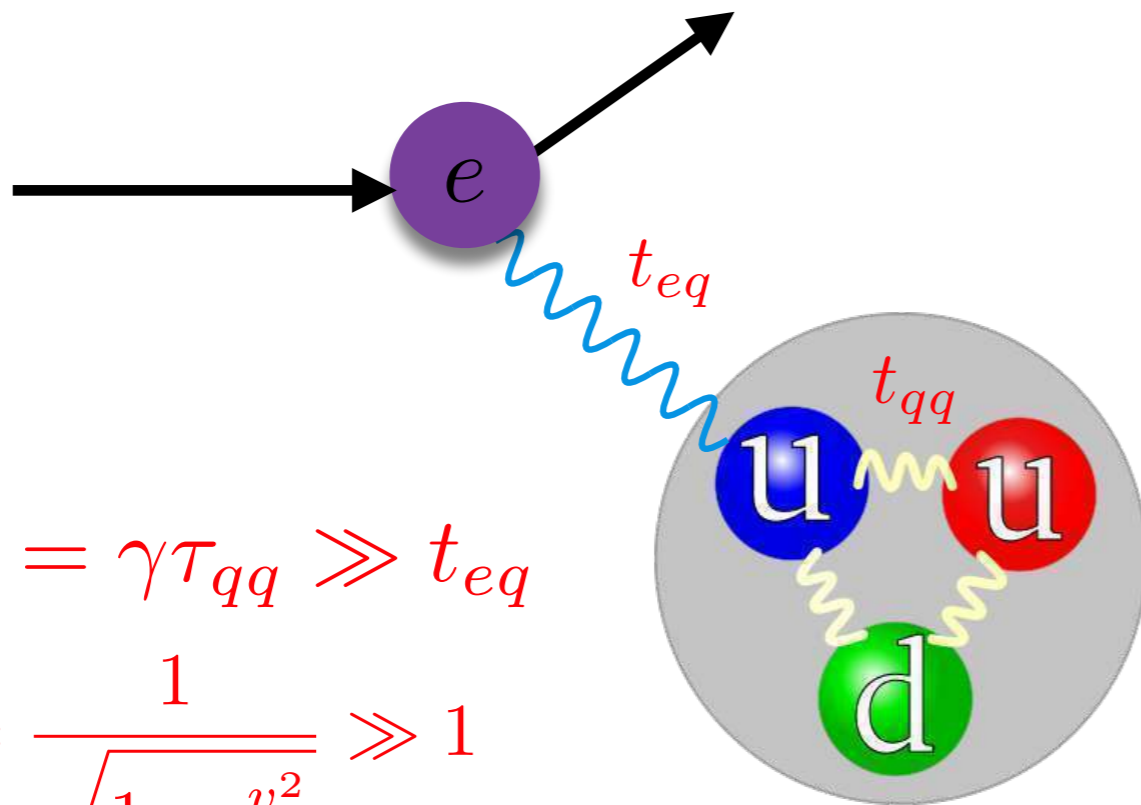
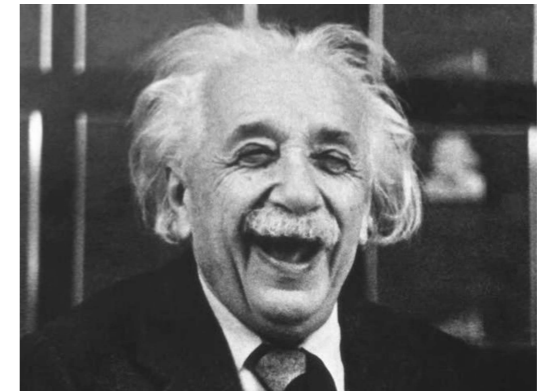
$$\sigma(ep) \sim \sigma(eq)$$



- But is this really sensible? Can we simply ignore fact that quarks is part of a complex and strongly bound system (the proton)?

Colliding Protons

- Proton at **rest**: complex system of interacting quarks ($\tau_{qq} \approx 10^{-24} s$).
- However we are interested in very **high energy** proton collisions. Proton has velocity $v \sim c$, and **relativity** comes into the game.
- What does electron 'see'? **Time dilation**: proton 'clock' much slower than when at rest, electron only sees a **static snapshot** of the proton!



$$t_{qq} = \gamma \tau_{qq} \gg t_{eq}$$

$$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} \gg 1$$

- Electron-quark interaction time \ll timescale of internal quark interactions!

Colliding Protons

- Electron scatters off a quark within the **static** proton **snapshot**. The **quark interactions** within the proton are **frozen** and can be ignored.
- Valid to consider in terms of **free** quark-electron scattering.

$$\sigma(ep) \sim \sigma(eq) \quad \checkmark$$

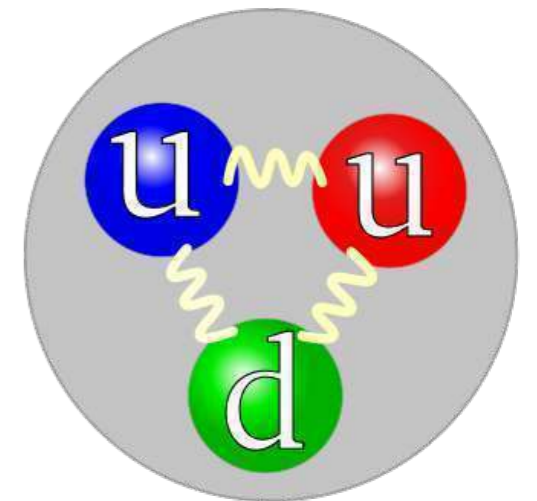
- Final element: what does the frozen distribution of quarks look like?
Relevant degree of freedom: amount of proton's energy carried by quark.

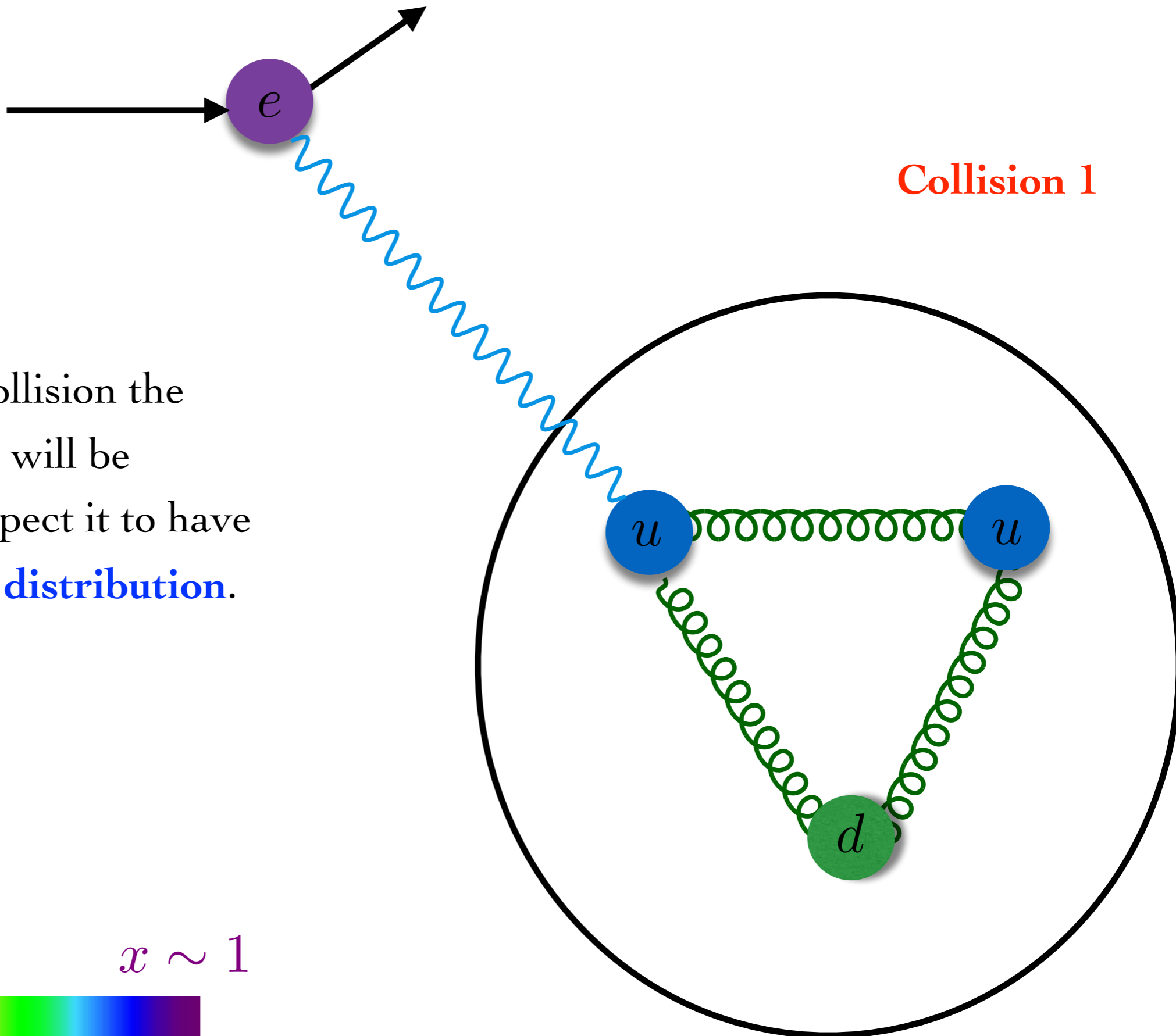
- Introduce **new variable**: $x = \frac{E_{\text{quark}}}{E_{\text{proton}}}$

$$0 < x < 1$$

↑ quark has no energy

↑ quark has all proton energy





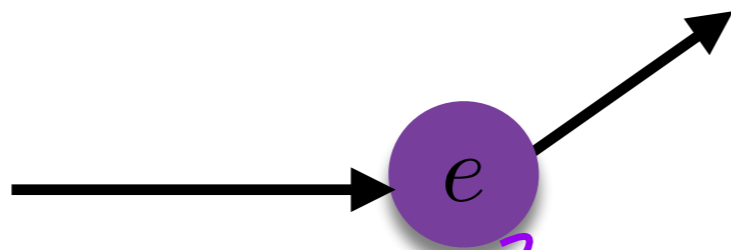
- Collision by collision the proton snapshot will be different, but expect it to have some **statistical distribution**.

$$x = \frac{E_{\text{quark}}}{E_{\text{proton}}}$$

$x \ll 1$

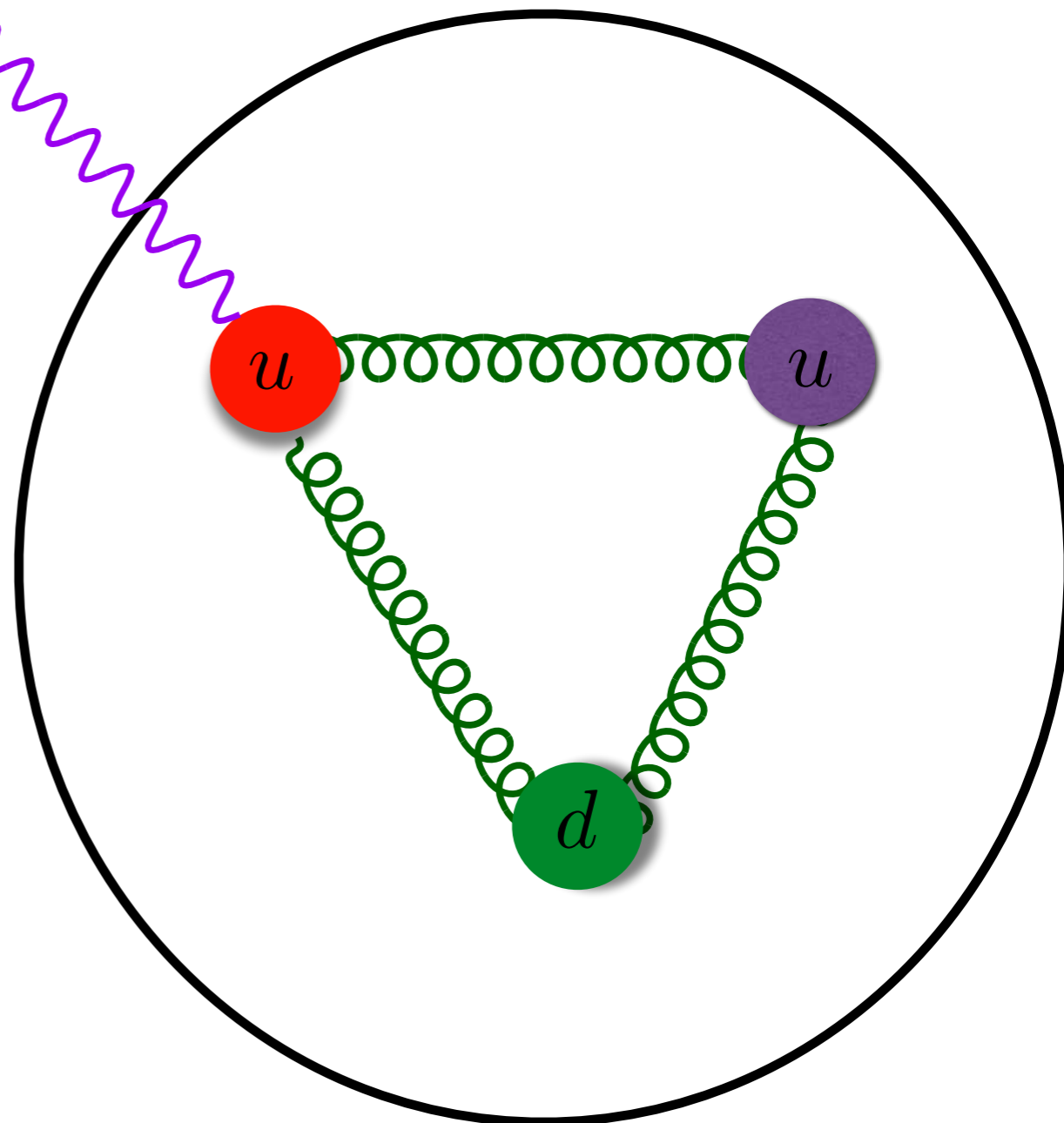
$x \sim 1$





Collision 2

- Collision by collision the proton snapshot will be different, but expect it to have some **statistical distribution**.

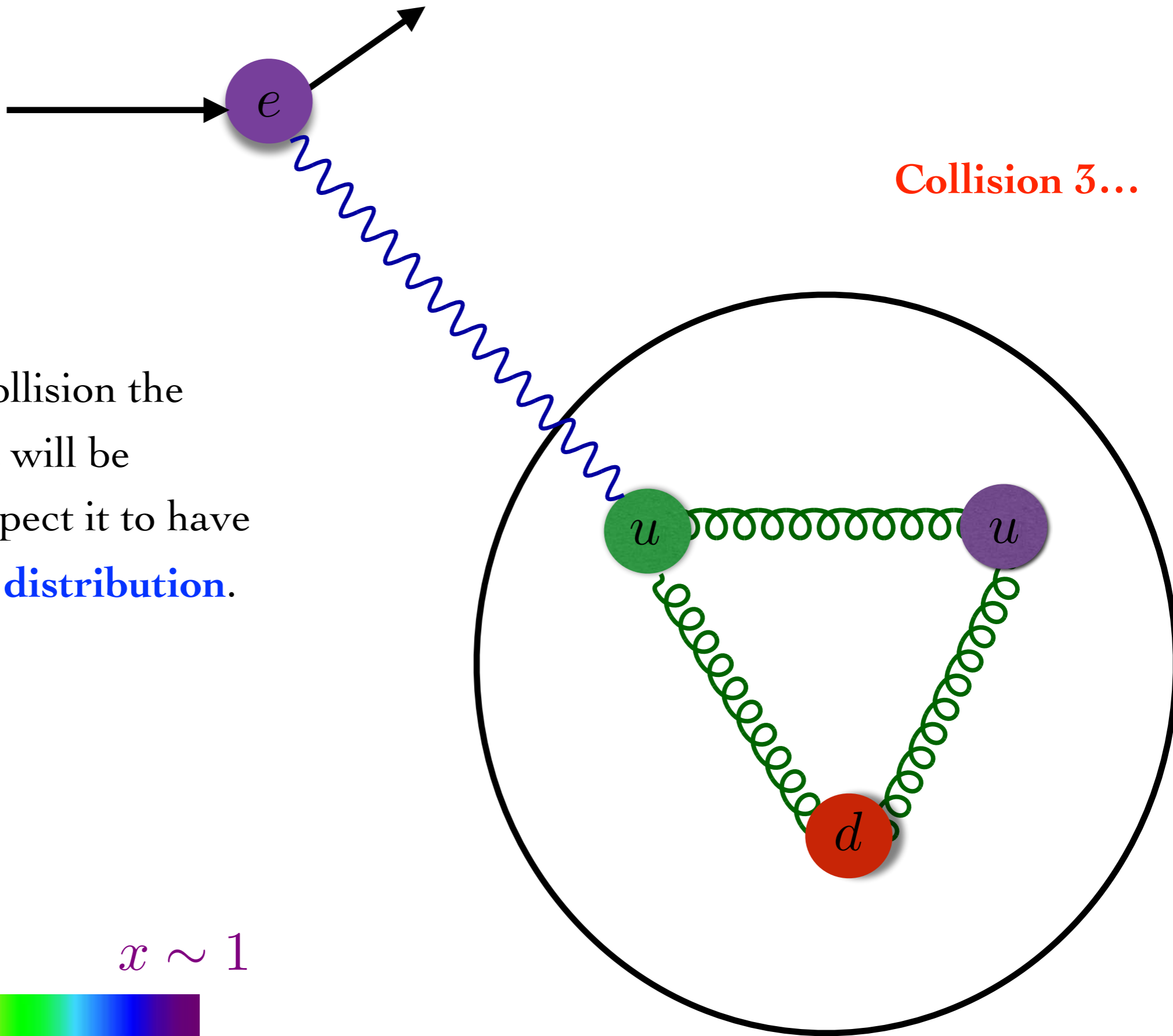


$$x = \frac{E_{\text{quark}}}{E_{\text{proton}}}$$

$x \ll 1$

$x \sim 1$





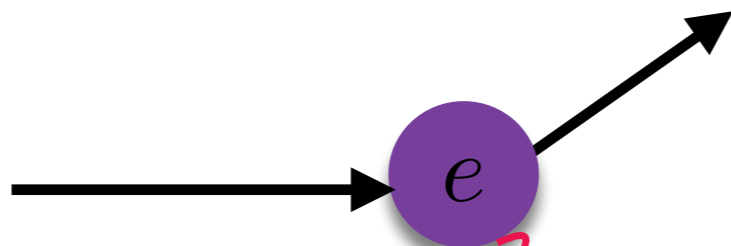
- Collision by collision the proton snapshot will be different, but expect it to have some **statistical distribution**.

$$x = \frac{E_{\text{quark}}}{E_{\text{proton}}}$$

$x \ll 1$

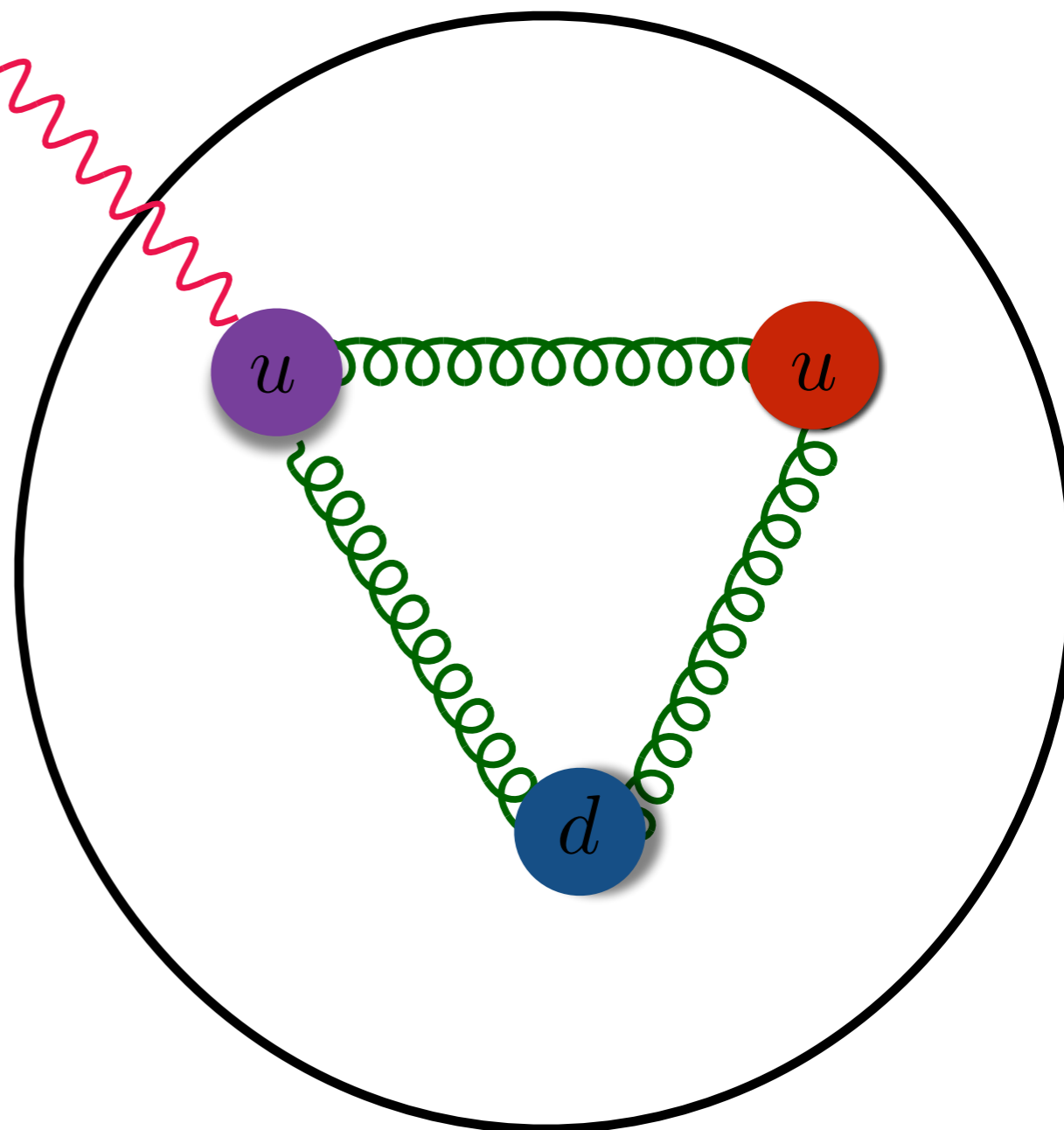
$x \sim 1$





...Collision N

- Collision by collision the proton snapshot will be different, but expect it to have some **statistical distribution**.



$$x = \frac{E_{\text{quark}}}{E_{\text{proton}}}$$

$x \ll 1$

$x \sim 1$



Mapping out the Proton

- **Statistical distribution** known as '**Parton Distribution Function**' (**PDF**).

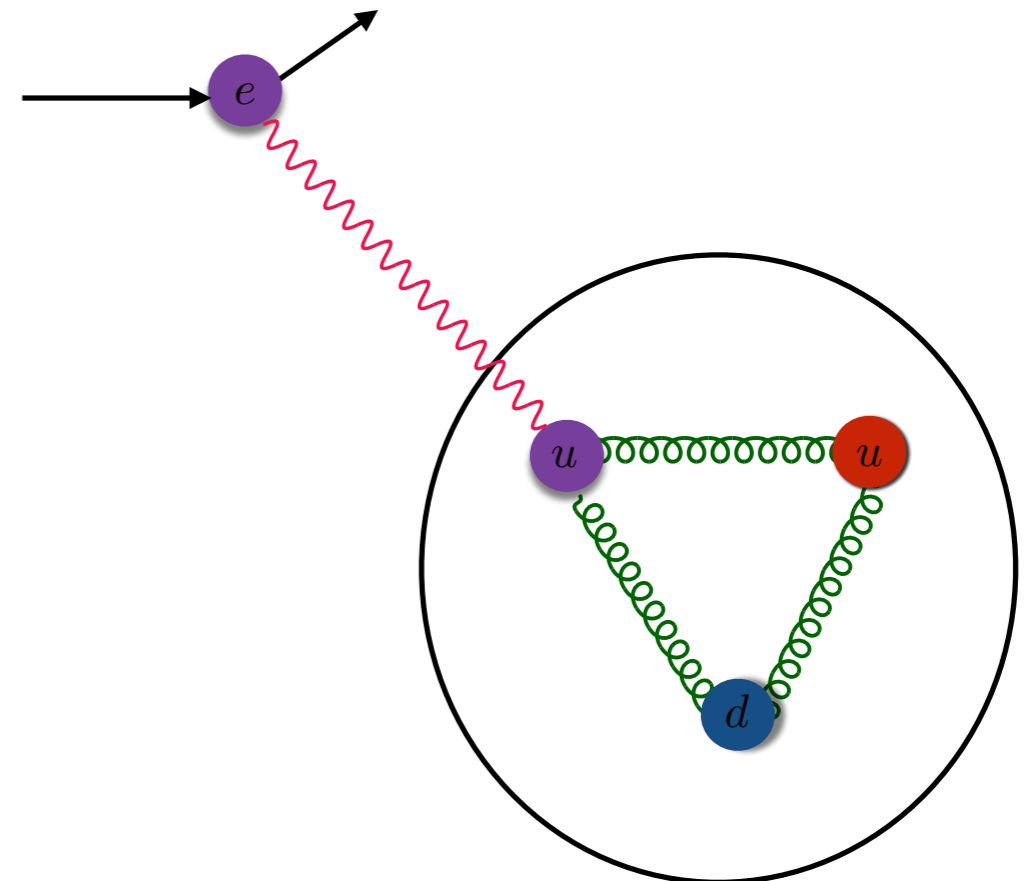
$f(x)$: **Probability of finding a quark with energy fraction x in proton snapshot.**

Electron-proton scattering rate:

$$\sigma_{ep} = \int_0^1 dx f_q(x) \sigma_{eq}(x)$$

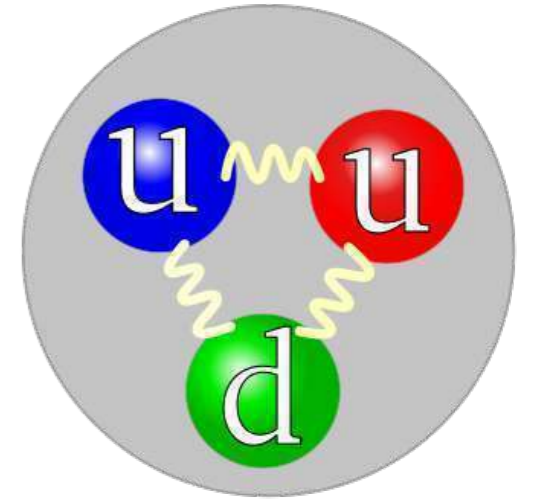
- Electron energy and scattering angle directly related to quark x and photon resolution.

→ By scanning over these, can **map out** entire distribution!



The Proton 'Sea'

- So far have only considered possibility of electron-quark scattering. Seems sensible given basic **uud** picture of proton. But is this the **whole story?**

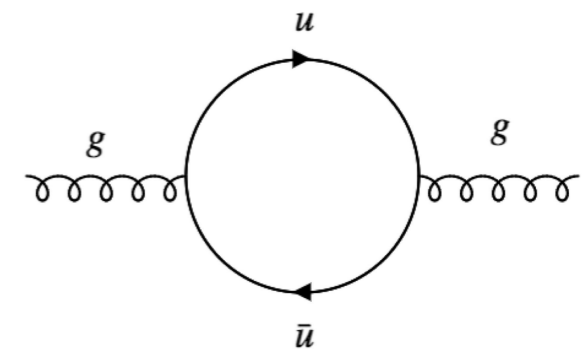


- **No!** In fact proton is a much more complex object than basic uud picture would predict. **Uncertainty principle:**

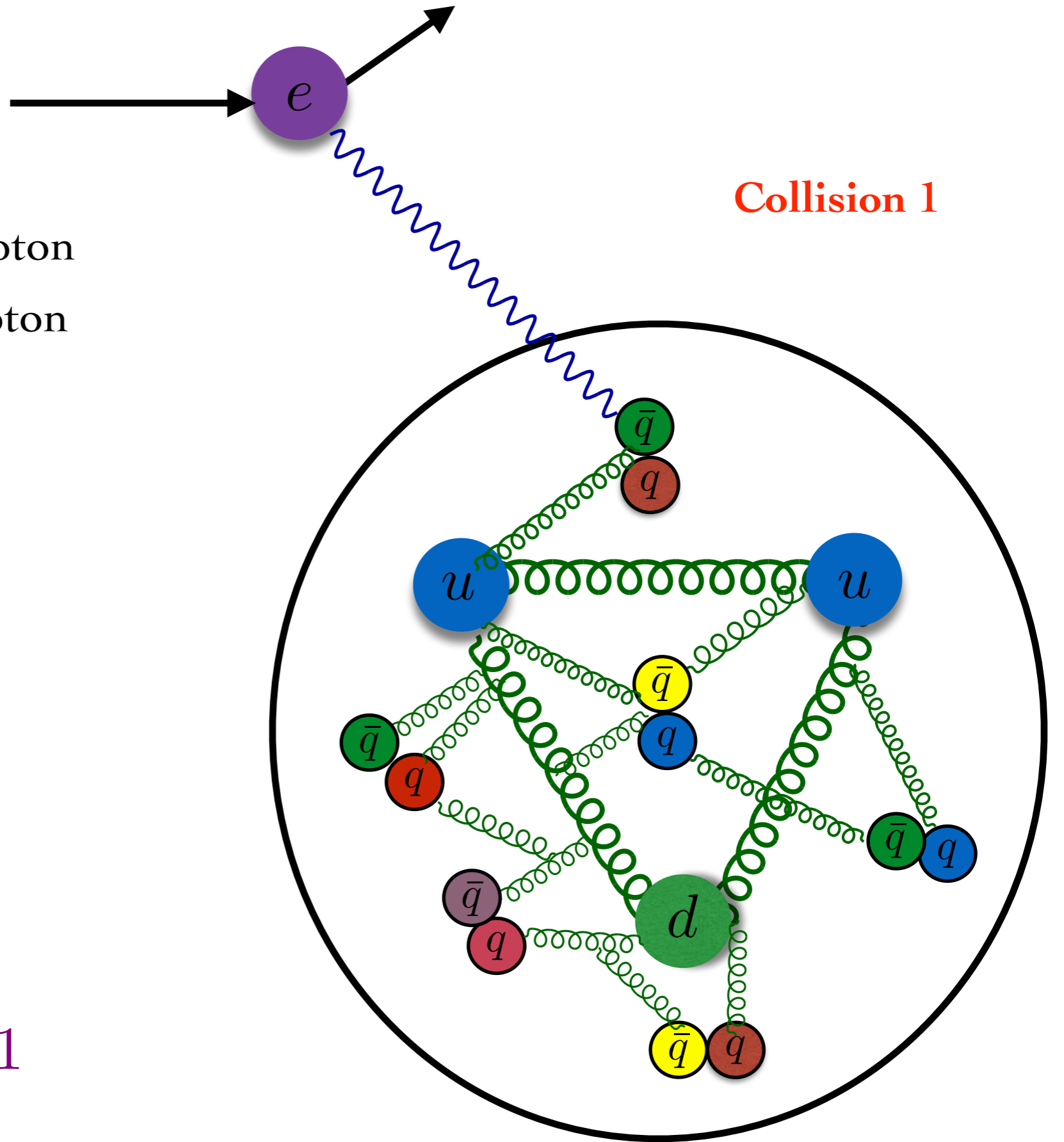
$$\Delta E \Delta t \sim \hbar$$

→ Quark-antiquark pairs can snap in and out of existence if:

$$\Delta t < \frac{\hbar}{2m_q c^2}$$



- Proton is in fact filled with a 'sea' of **quark-antiquark pairs** snapping in and out of existence. In addition to so-called 'valence' **uud**.



Collision 1

- Back to the electron-proton collision: this seething proton sea is also ~ **frozen** in the proton snapshot.

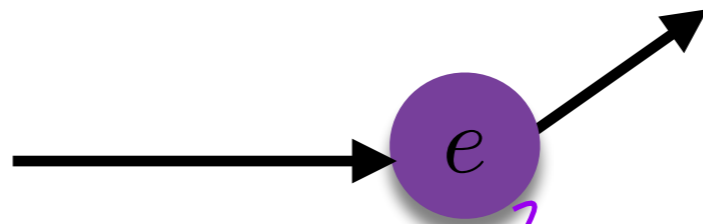
→ Electron can scatter off these frozen quark-antiquark pairs!

$$x = \frac{E_{\text{quark}}}{E_{\text{proton}}}$$

$x \ll 1$

$x \sim 1$

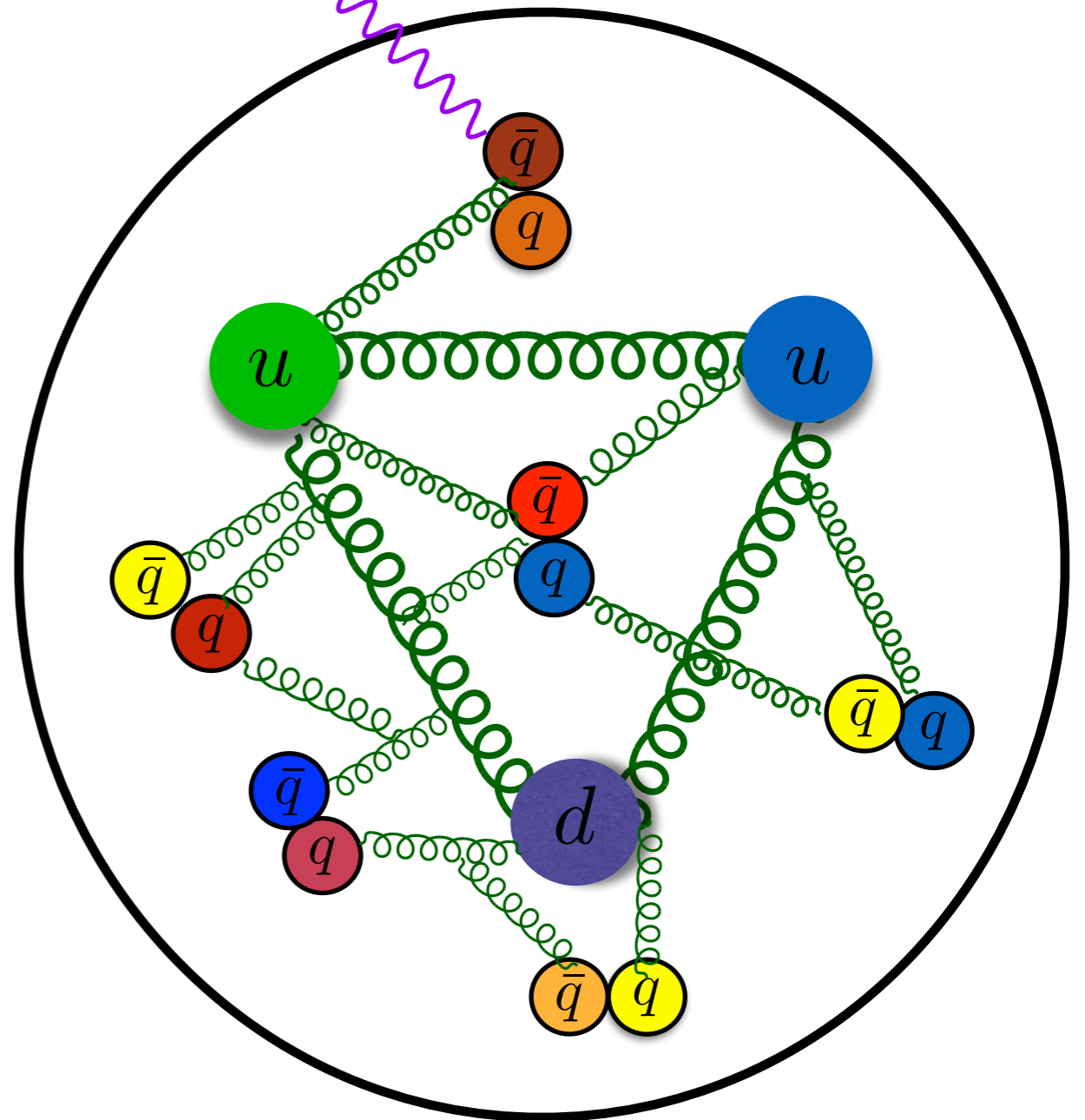




Collision 2...

- Back to the electron-proton collision: this seething proton sea is also ~ **frozen** in the proton snapshot.

→ Electron can scatter off these frozen quark-antiquark pairs!

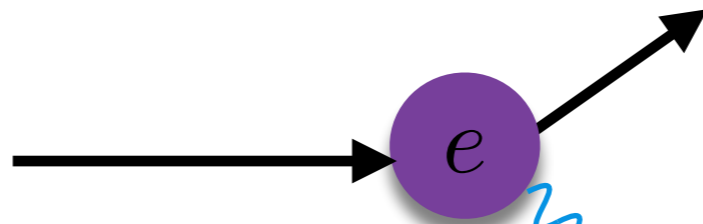


$$x = \frac{E_{\text{quark}}}{E_{\text{proton}}}$$

$x \ll 1$

$x \sim 1$





...Collision N

- Scattering rate becomes:

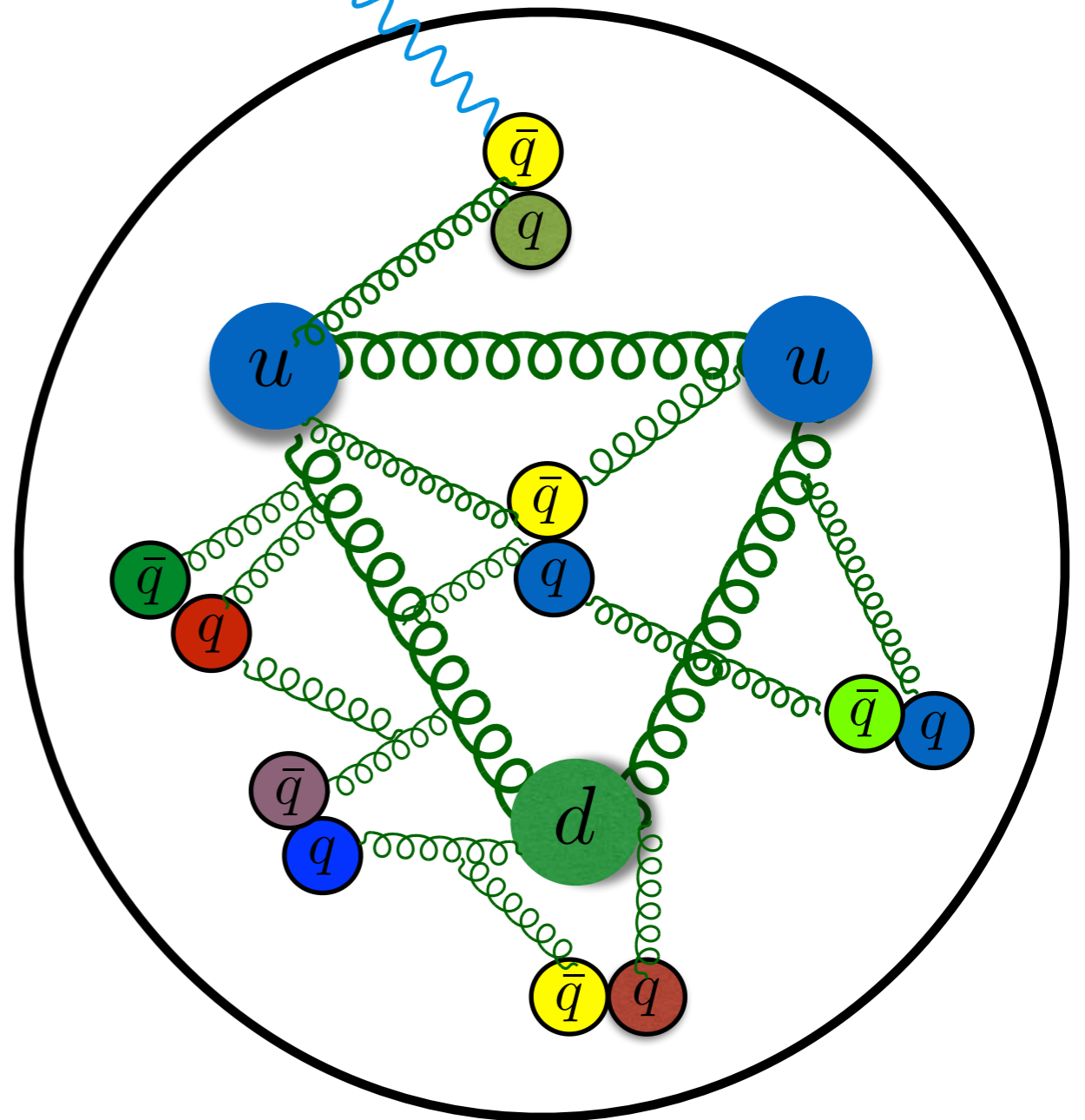
$$\sigma_{ep} = \int_0^1 dx [f_q(x) \sigma_{eq}(x) + f_{\bar{q}}(x) \sigma_{e\bar{q}}(x)]$$

$f_{\bar{q}}(x)$: Probability to find **antiquark** with energy fraction x within proton snapshot.

$$x = \frac{E_{\text{quark}}}{E_{\text{proton}}}$$

$x \ll 1$

$x \sim 1$



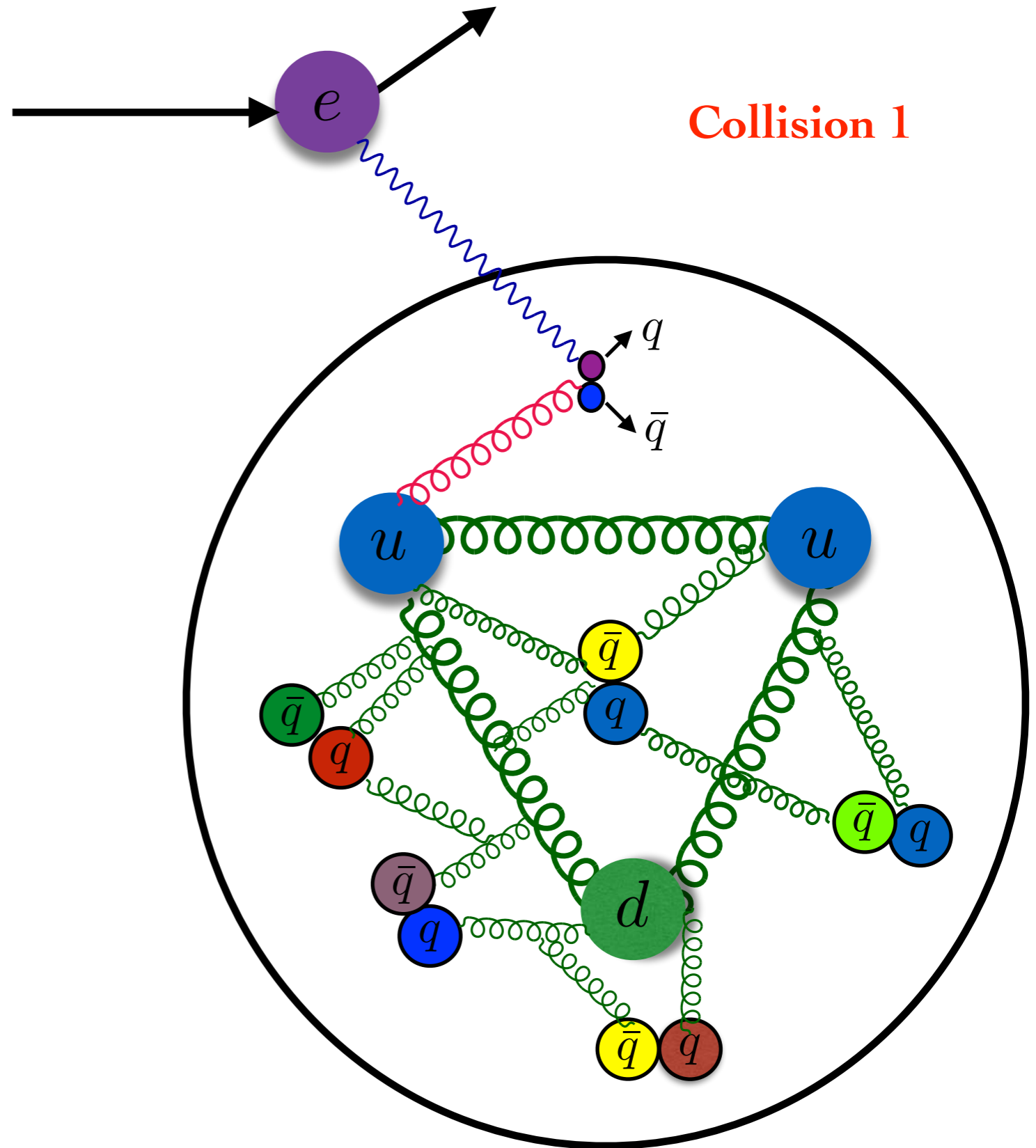
The Gluon Content of the Proton

- Sea of **gluons** constantly being exchanged between the quark constituents of proton, binding it together.
- Can also scatter off these objects!

$$x = \frac{E_{\text{quark}}}{E_{\text{proton}}}$$

$$x \ll 1$$

$$x \sim 1$$

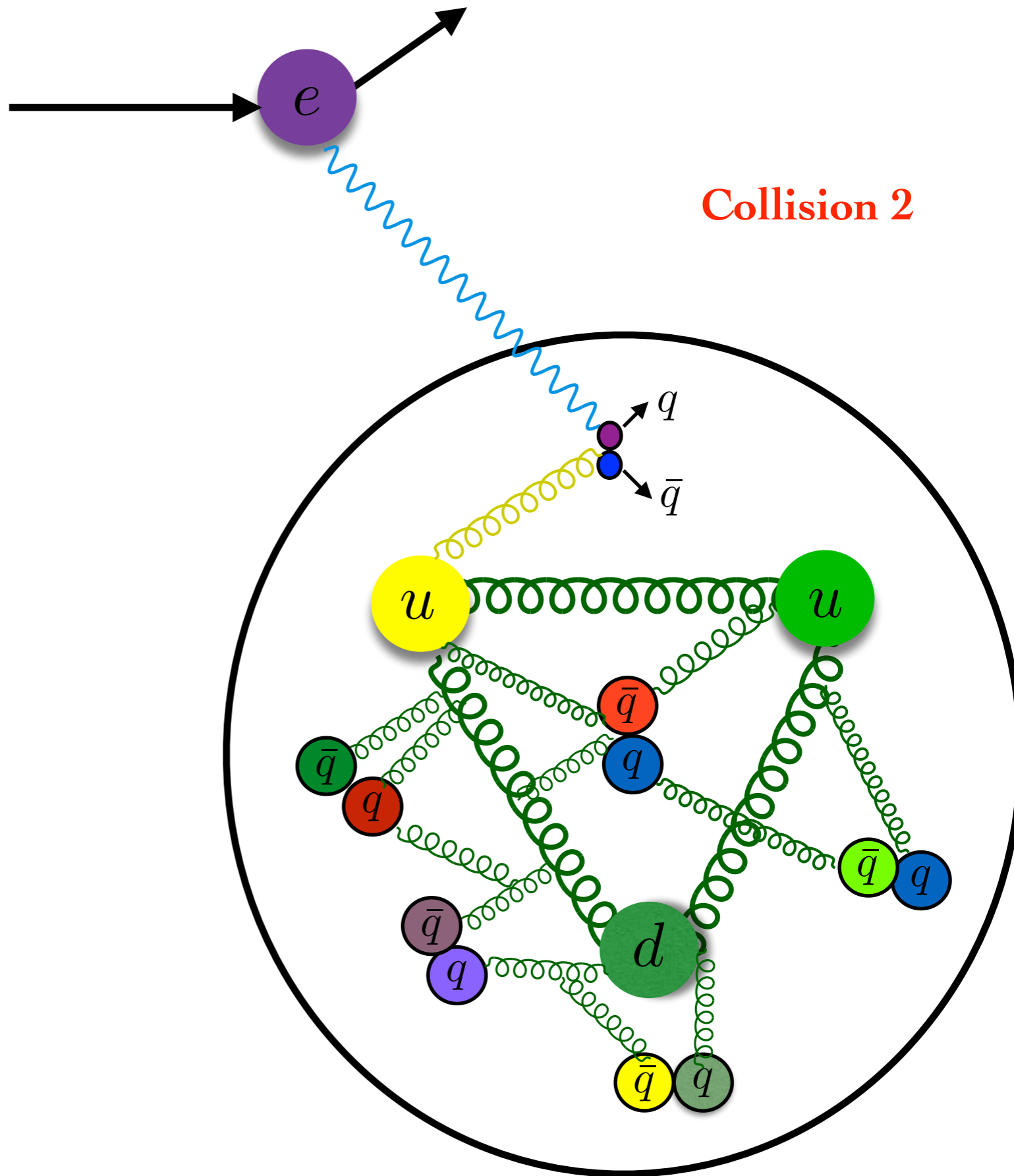


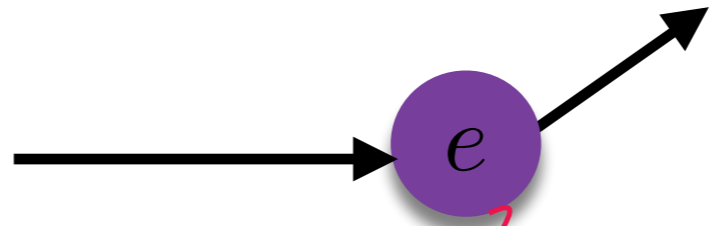
- Sea of **gluons** constantly being exchanged between the quark constituents of proton, binding it together.
- Can also scatter off these objects!

$$x = \frac{E_{\text{quark}}}{E_{\text{proton}}}$$

$x \ll 1$

$x \sim 1$





...Collision N

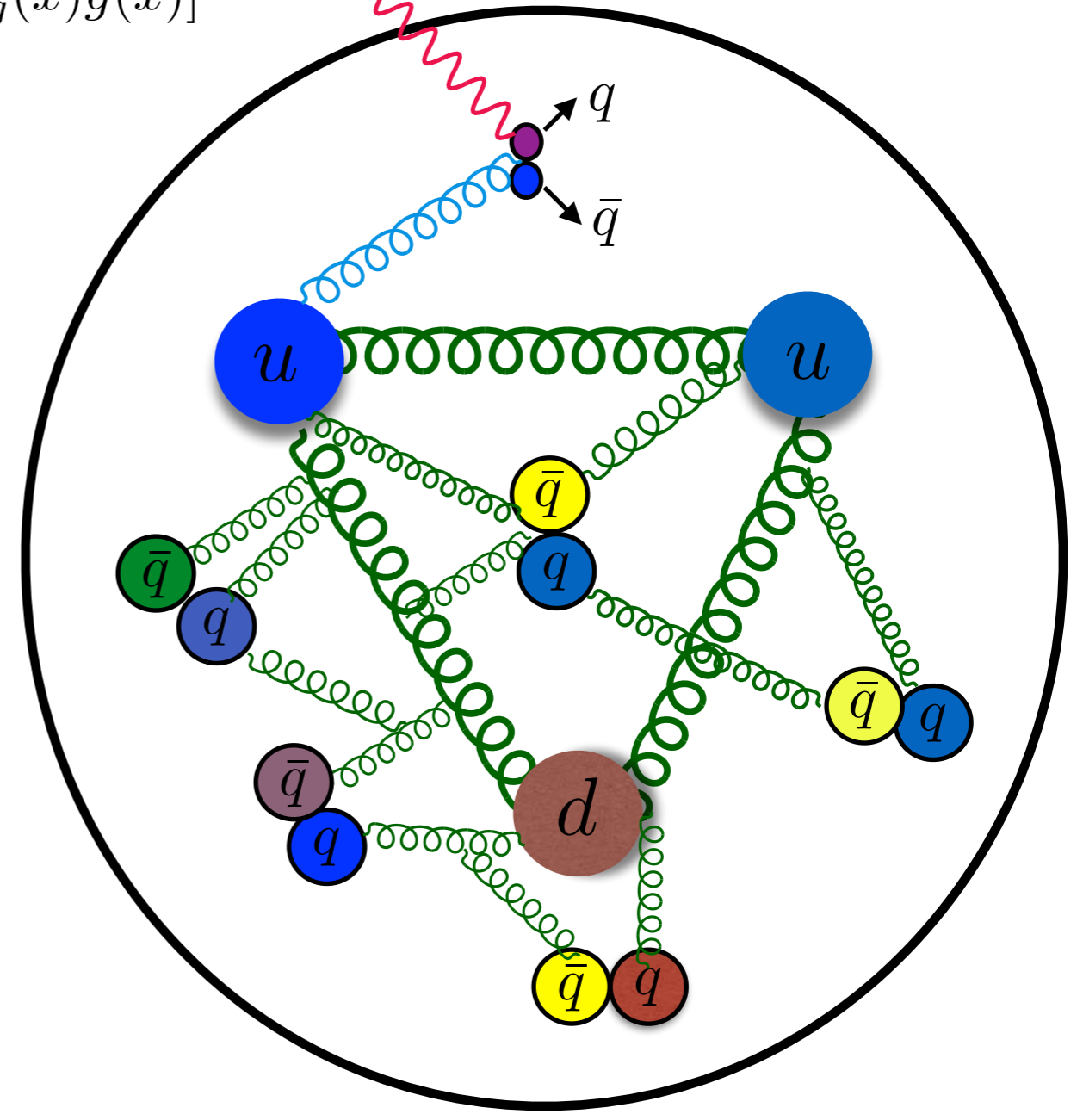
$$\sigma_{ep} = \int_0^1 dx [f_q(x) \sigma_{eq}(x) + f_{\bar{q}}(x) \sigma_{e\bar{q}}(x) + \sigma_{eg}(x)g(x)]$$

$g(x)$: Probability to find gluon with energy fraction x within proton snapshot.

$$x = \frac{E_{\text{quark}}}{E_{\text{proton}}}$$

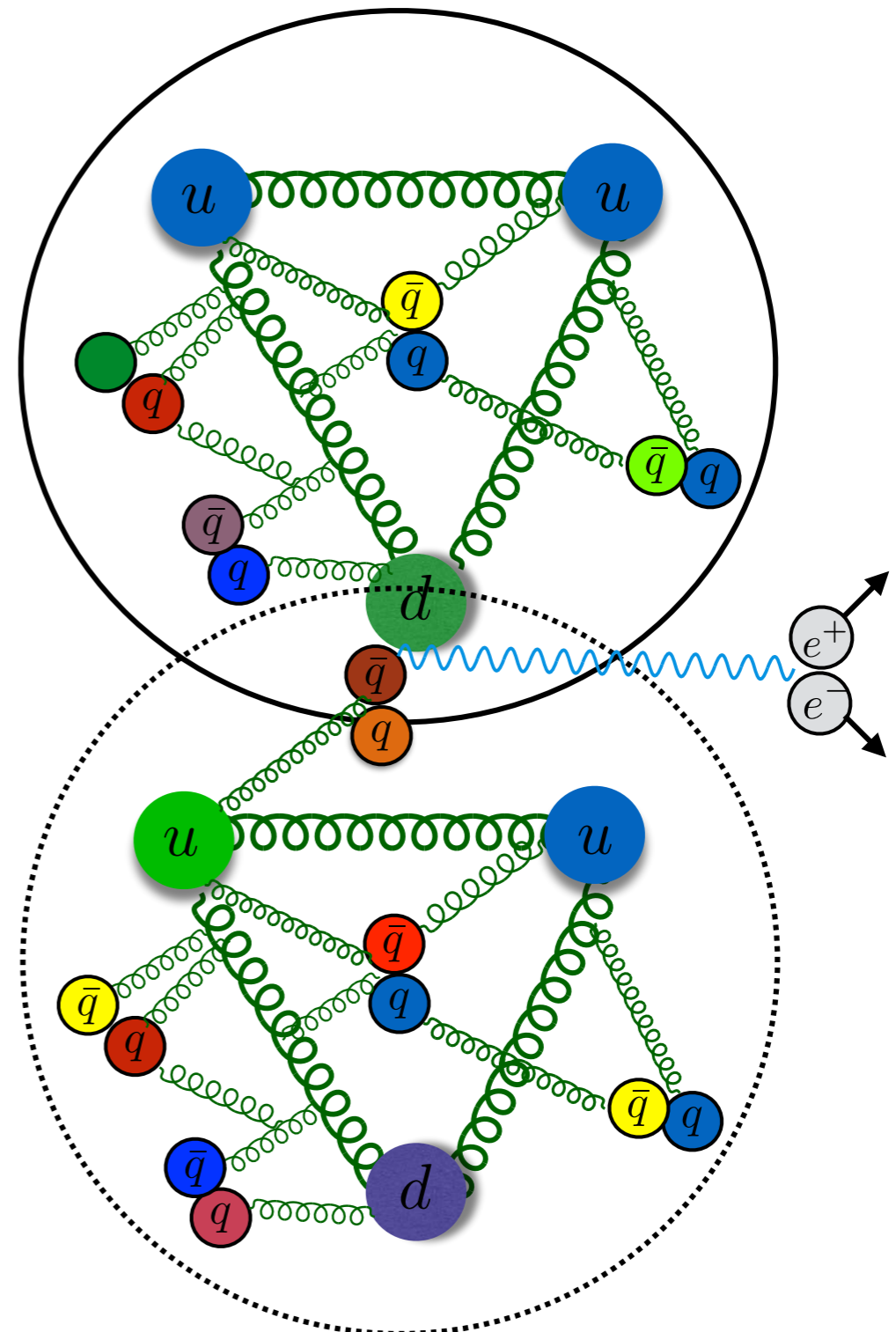
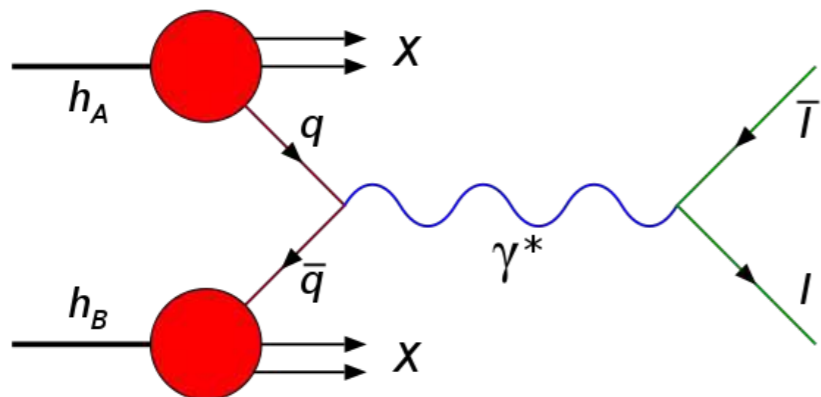
$x \ll 1$

$x \sim 1$



Proton-Proton Collisions

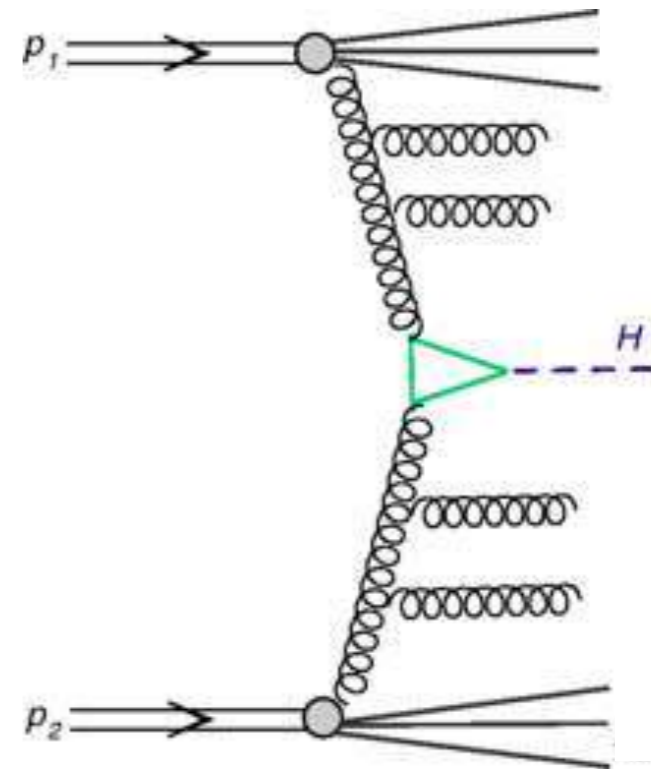
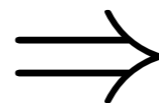
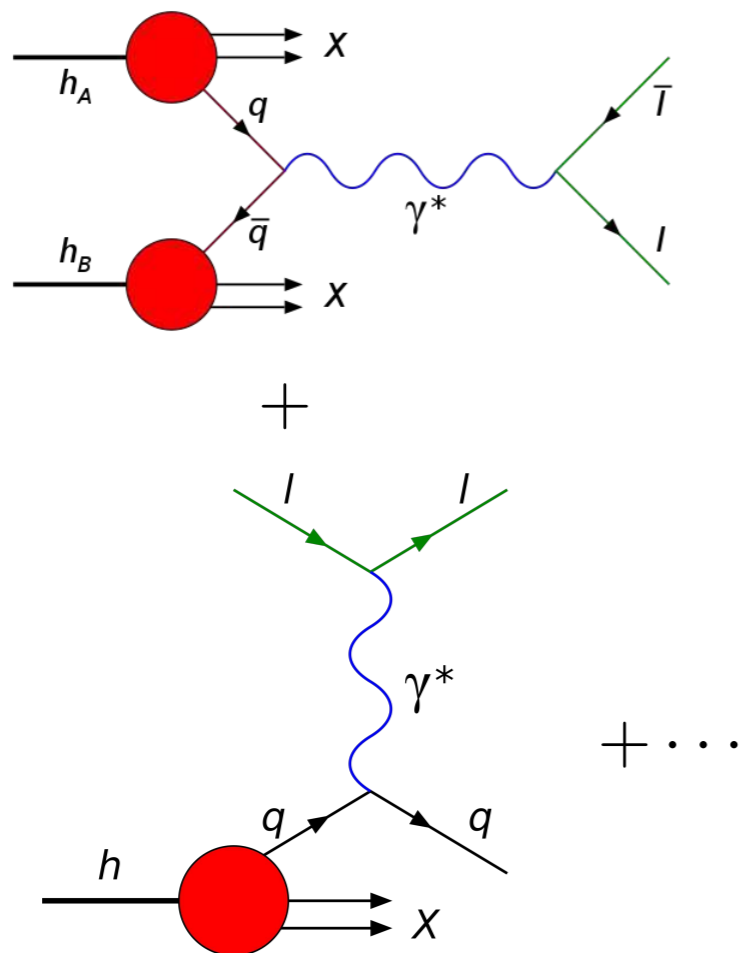
- How does this picture change in the case of **proton-proton** collisions at the **LHC**? Answer: **not too much**.
- Proton-proton scattering = scattering of point-like **quark/gluons**.
- For example, simple so-called '**Drell-Yan**' process, leading to production of an electron-positron pair.



$$\sigma^{DY} \sim \sigma(\bar{q}q \rightarrow e^+e^-) \otimes q(x_1) \otimes \bar{q}(x_2)$$

Extracting PDFs

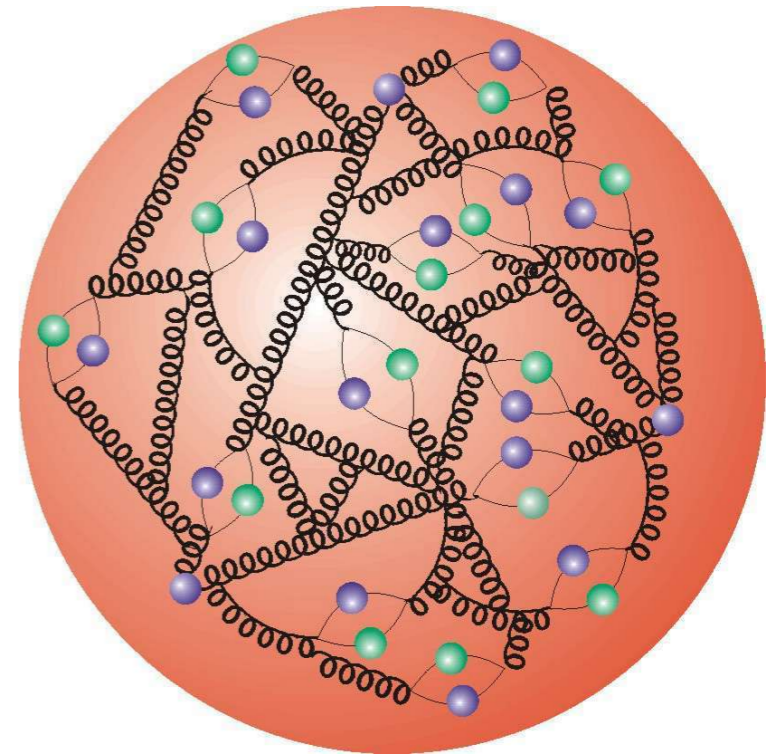
- How do we actually **determine** these **PDFs**?
- **PDFs**: distribution of proton constituents, due to complex strong interactions. With current theoretical tools we **cannot predict** them.
- What we can do: determine them experimentally in processes we understand well and use to predict processes such as **Higgs** production.



The Proton Today

- Proton: a complex quantum-mechanical system. Active and ongoing programme of research to **map** it out:

- ❖ How much of proton energy is carried by **gluons**? How much is in the quark **sea**?
- ❖ How significant are **heavier** (strange, charm, bottom, top...) **quarks**?
- ❖ How do each of these vary with x ?
- ❖ How well do we understand the connection between the proton and high-energy collisions?



- To probe the Higgs and stress test the Standard Model at the LHC we need to address these questions as **precisely** as we can.

Global PDF Fits

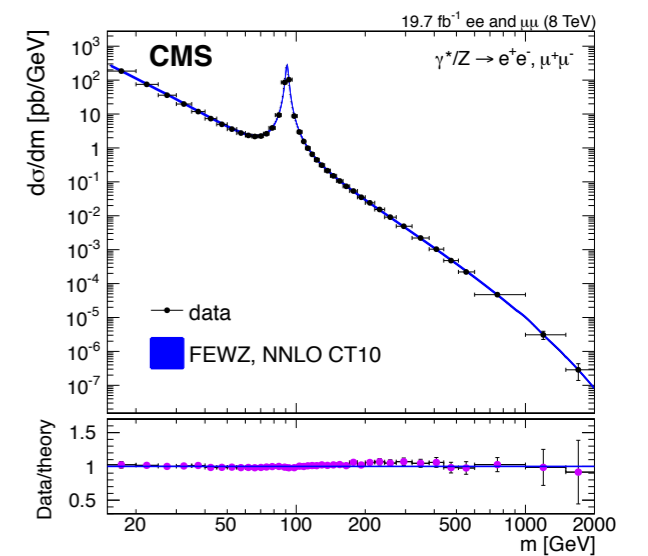
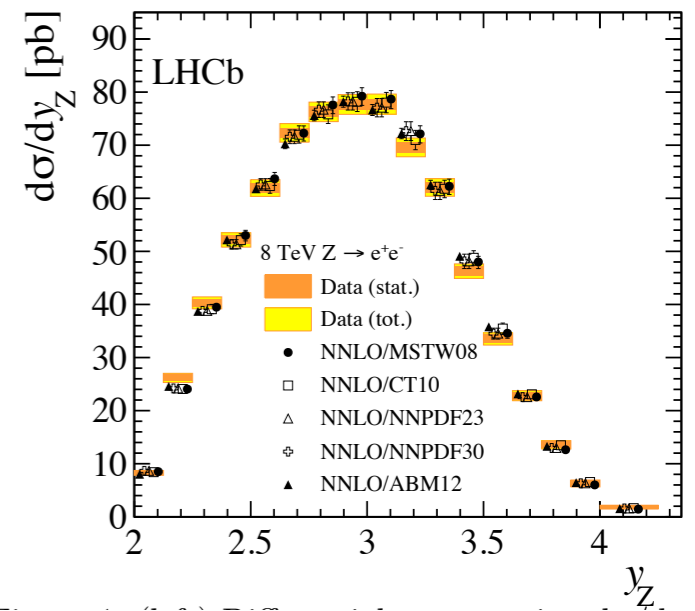
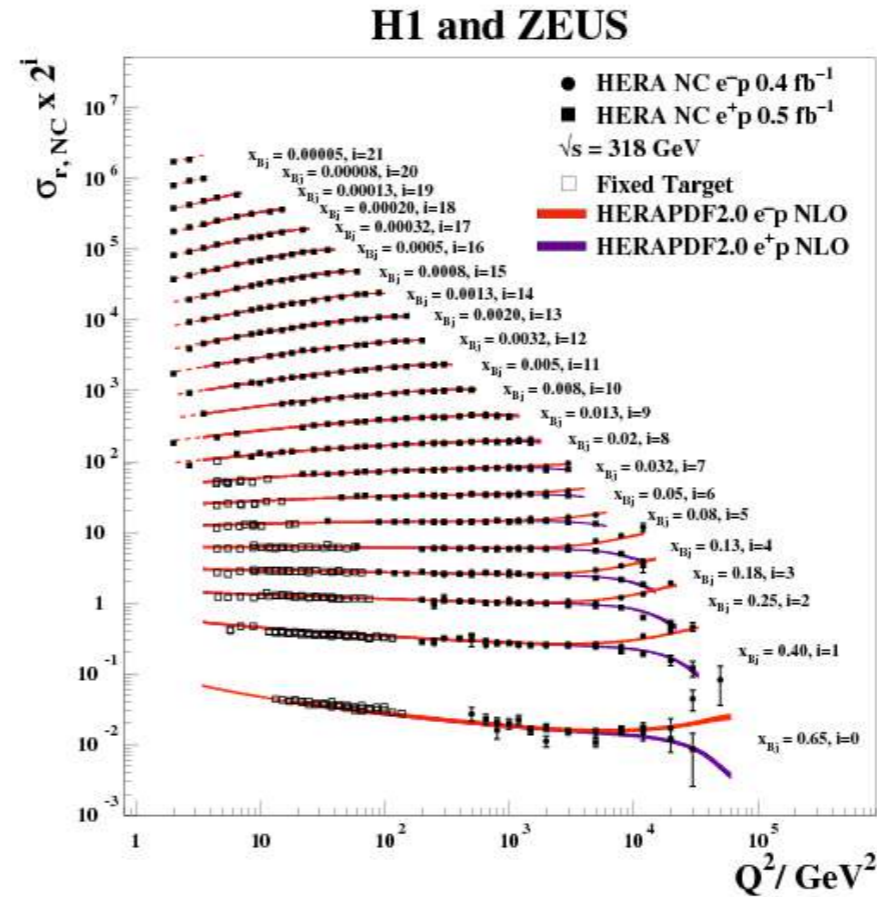
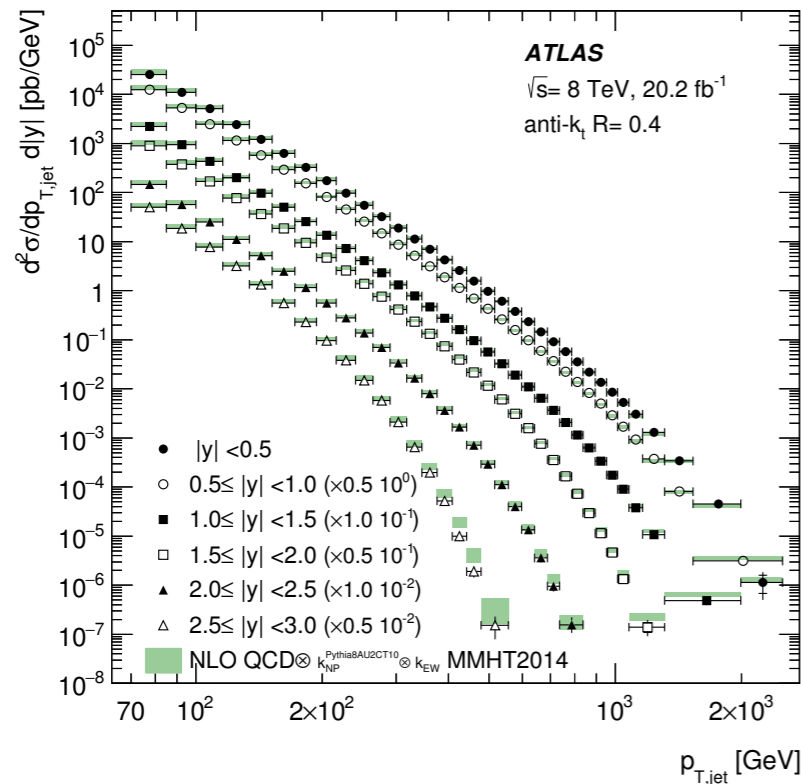
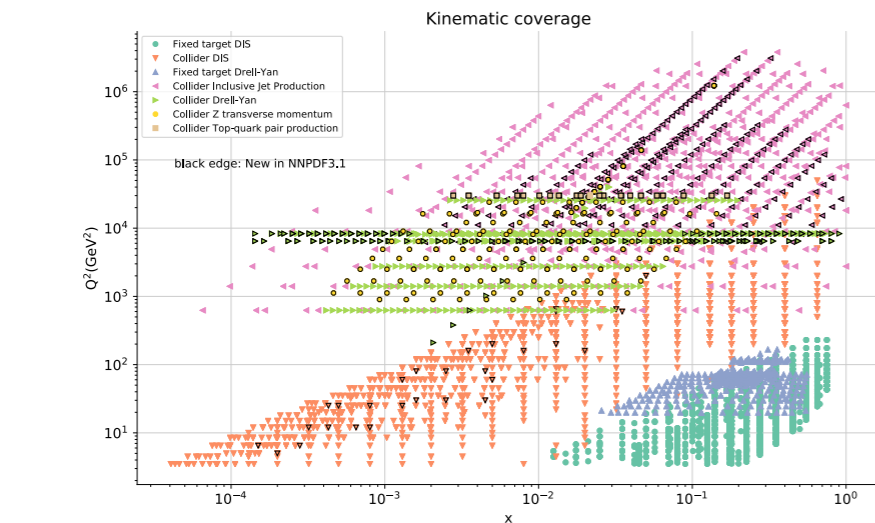
- Since the discovery of quarks, a vast array of data on proton collisions has been collected, from different **colliders** and via different **processes**.
- Combining all of this we can perform a ‘**Global Fit**’ to PDFs, allowing us to pin down proton structure with **precision**.

Data set	LO	NLO	NNLO
BCDMS $\mu p F_2$ [125]	162 / 153	176 / 163	173 / 163
BCDMS $\mu d F_2$ [19]	140 / 142	143 / 151	143 / 151
NMC $\mu p F_2$ [20]	141 / 115	132 / 123	123 / 123
NMC $\mu d F_2$ [20]	134 / 115	115 / 123	108 / 123
NMC $\mu n / \mu p$ [21]	122 / 137	131 / 148	127 / 148
E665 $\mu p F_2$ [22]	59 / 53	60 / 53	65 / 53
E665 $\mu d F_2$ [22]	52 / 53	52 / 53	60 / 53
SLAC $ep F_2$ [23, 24]	21 / 18	31 / 37	31 / 37
SLAC $ed F_2$ [23, 24]	13 / 18	30 / 38	26 / 38
NMC/BCDMS/SLAC/HERA F_L [20, 125, 24, 63, 64, 65]	113 / 53	68 / 57	63 / 57
ES66/NuSea pp DY [88]	229 / 184	221 / 184	227 / 184
ES66/NuSea pd/pp DY [89]	29 / 15	11 / 15	11 / 15
NuTeV $\nu N F_2$ [29]	35 / 49	39 / 53	38 / 53
CHORUS $\nu N F_2$ [30]	25 / 37	26 / 42	28 / 42
NuTeV $\nu N xF_3$ [29]	49 / 42	37 / 42	31 / 42
CHORUS $\nu N xF_3$ [30]	35 / 28	22 / 28	19 / 28
CCFR $\nu N \rightarrow \mu\mu X$ [31]	65 / 86	71 / 86	76 / 86
NuTeV $\nu N \rightarrow \mu\mu X$ [31]	53 / 40	38 / 40	43 / 40
HERA e^+p NC 820 GeV [61]	125 / 78	93 / 78	89 / 78
HERA e^+p NC 920 GeV [61]	479 / 330	402 / 330	373 / 330
HERA e^-p NC 920 GeV [61]	158 / 145	129 / 145	125 / 145
HERA e^+p CC [61]	41 / 34	34 / 34	32 / 34
HERA e^-p CC [61]	29 / 34	23 / 34	21 / 34
HERA $ep F_2^{\text{charm}}$ [62]	105 / 52	72 / 52	82 / 52
H1 99-00 e^+p incl. jets [126]	77 / 24	14 / 24	—
ZEUS incl. jets [127, 128]	140 / 60	45 / 60	—
DØ II pp incl. jets [119]	125 / 110	116 / 110	119 / 110
CDF II pp incl. jets [118]	78 / 76	63 / 76	59 / 76
CDF II W asym. [66]	55 / 13	32 / 13	30 / 13
DØ II $W \rightarrow \nu e$ asym. [67]	47 / 12	28 / 12	27 / 12
DØ II $W \rightarrow \nu \mu$ asym. [68]	16 / 10	19 / 10	21 / 10
DØ II Z rap. [90]	34 / 28	16 / 28	16 / 28
CDF II Z rap. [70]	95 / 28	36 / 28	40 / 28
ATLAS W^+, W^-, Z [10]	94/30	38/30	39/30
CMS W asymm $p_T > 35$ GeV [9]	10/11	7/11	9/11
CMS asymm $p_T > 25$ GeV, 30 GeV [77]	7/24	8/24	10/24
LHCb $Z \rightarrow e^+e^-$ [79]	76/9	13/9	20/9
LHCb W asymm $p_T > 20$ GeV [78]	27/10	12/10	16/10
CMS $Z \rightarrow e^+e^-$ [84]	46/35	19/35	22/35
ATLAS high-mass Drell-Yan [83]	42/13	21/13	17/13
CMS double diff. Drell-Yan [86]	—	372/132	149/132
Tevatron, ATLAS, CMS σ_H [91]-[97]	53/13	7/13	8/13
ATLAS jets (2.76 TeV+7 TeV) [108, 107]	162/116	106/116	—
CMS jets (7 TeV) [106]	150/133	138/133	—



Where are we now?

- The wealth of available high precision data is unprecedented.



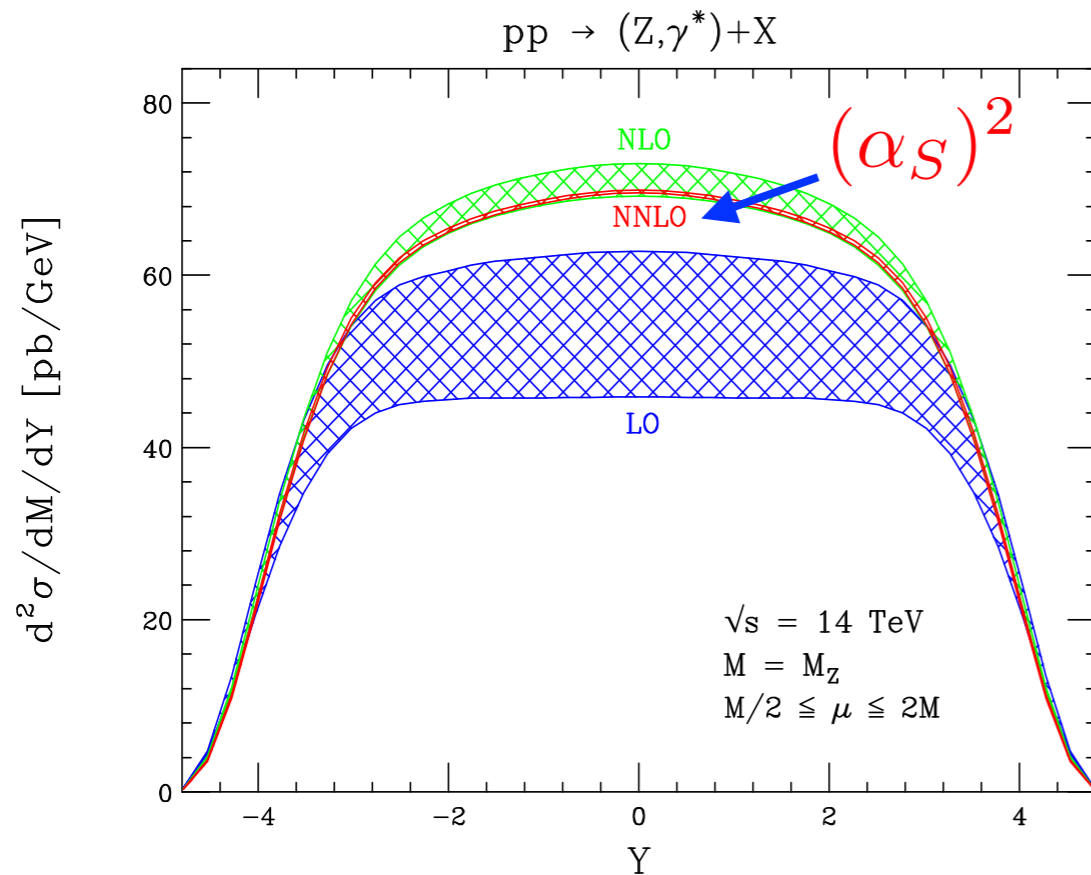
- There is **a lot!** And much of it these days comes from **LHC** itself.

Where are we now?

- The **theoretical calculations** for the processes entering PDF fits are certainly keeping up with the data!

Five-loop contributions to low- N non-singlet anomalous dimensions in QCD

$(\alpha_s)^5$



$$\gamma_{\text{ns}}^{(4)+}(N=2) =$$

$$C_F^5 \left[\frac{9306376}{19683} - \frac{802784}{729} \zeta_3 - \frac{557440}{81} \zeta_5 + \frac{12544}{9} \zeta_3^2 + 8512 \zeta_7 \right]$$

$$- C_A C_F^4 \left[\frac{81862744}{19683} - \frac{1600592}{243} \zeta_3 + \frac{59840}{81} \zeta_4 - \frac{142240}{27} \zeta_5 + 3072 \zeta_3^2 - \frac{35200}{9} \zeta_6 + 19936 \zeta_7 \right]$$

$$+ C_A^2 C_F^3 \left[\frac{63340406}{6561} - \frac{1003192}{243} \zeta_3 - \frac{229472}{81} \zeta_4 + \frac{61696}{27} \zeta_5 + \frac{30976}{9} \zeta_3^2 - \frac{35200}{9} \zeta_6 + 15680 \zeta_7 \right]$$

$$- C_A^3 C_F^2 \left[\frac{220224724}{19683} + \frac{4115536}{729} \zeta_3 - \frac{170968}{27} \zeta_4 - \frac{3640624}{243} \zeta_5 + \frac{70400}{27} \zeta_3^2 + \frac{123200}{27} \zeta_6 + \frac{331856}{27} \zeta_7 \right]$$

$$+ C_A^4 C_F \left[\frac{266532611}{39366} + \frac{2588144}{729} \zeta_3 - \frac{221920}{81} \zeta_4 - \frac{3102208}{243} \zeta_5 + \frac{74912}{81} \zeta_3^2 + \frac{334400}{81} \zeta_6 + \frac{178976}{27} \zeta_7 \right]$$

$$- \frac{d_{NA}^{(4)}}{N_A} C_F \left[\frac{15344}{81} - \frac{12064}{27} \zeta_3 - \frac{704}{3} \zeta_4 + \frac{58400}{27} \zeta_5 - \frac{6016}{3} \zeta_3^2 - \frac{19040}{9} \zeta_7 \right]$$

$$+ \frac{d_{NF}^{(4)}}{N_F} C_F \left[\frac{23968}{81} - \frac{733504}{81} \zeta_3 + \frac{176320}{81} \zeta_5 + \frac{6400}{3} \zeta_3^2 + \frac{77056}{9} \zeta_7 \right]$$

$$- \frac{d_{NF}^{(4)}}{N_F} C_A \left[\frac{82768}{81} - \frac{555520}{81} \zeta_3 + \frac{10912}{9} \zeta_4 - \frac{1292960}{81} \zeta_5 + \frac{84352}{27} \zeta_3^2 + \frac{140800}{27} \zeta_6 + 12768 \zeta_7 \right]$$

$$+ n_f C_F^4 \left[\frac{1824964}{19683} - \frac{463520}{243} \zeta_3 + \frac{21248}{81} \zeta_4 - \frac{16480}{81} \zeta_5 + \frac{6656}{9} \zeta_3^2 - \frac{6400}{9} \zeta_6 + \frac{8960}{3} \zeta_7 \right]$$

$$- n_f C_A C_F^3 \left[\frac{3375082}{6561} - \frac{420068}{243} \zeta_3 - \frac{48256}{81} \zeta_4 + \frac{458032}{81} \zeta_5 + \frac{3968}{3} \zeta_3^2 - \frac{8000}{3} \zeta_6 + \frac{4480}{3} \zeta_7 \right]$$

$$+ n_f C_A^2 C_F^2 \left[\frac{15291499}{13122} + \frac{1561600}{243} \zeta_3 - \frac{114536}{27} \zeta_4 - \frac{252544}{243} \zeta_5 + \frac{24896}{27} \zeta_3^2 + \frac{13600}{27} \zeta_6 + \frac{11200}{27} \zeta_7 \right]$$

$$- n_f C_A^3 C_F \left[\frac{48846580}{19683} + \frac{4314308}{729} \zeta_3 - \frac{274768}{81} \zeta_4 - \frac{1389080}{243} \zeta_5 + \frac{27808}{81} \zeta_3^2 + \frac{184000}{81} \zeta_6 + \frac{3908}{27} \zeta_7 \right]$$

$$+ n_f \frac{d_{NF}^{(4)}}{N_F} \left[\frac{22096}{27} + \frac{43712}{81} \zeta_3 - \frac{512}{9} \zeta_4 - \frac{217280}{81} \zeta_5 - \frac{25088}{27} \zeta_3^2 + \frac{25600}{27} \zeta_6 - 2464 \zeta_7 \right]$$

$$- n_f C_F \frac{d_{FF}^{(4)}}{N_F} \left[\frac{170752}{81} - \frac{328832}{81} \zeta_3 + \frac{650240}{81} \zeta_5 - \frac{8192}{9} \zeta_3^2 - \frac{35840}{9} \zeta_7 \right]$$

$$+ n_f C_A \frac{d_{FF}^{(4)}}{N_F} \left[\frac{207824}{81} + \frac{251392}{81} \zeta_3 - \frac{5632}{9} \zeta_4 - \frac{522880}{81} \zeta_5 + \frac{15872}{27} \zeta_3^2 + \frac{70400}{27} \zeta_6 - \frac{29120}{9} \zeta_7 \right]$$

$$+ n_f^2 C_F^3 \left[\frac{1082297}{6561} - \frac{145792}{243} \zeta_3 + \frac{1072}{81} \zeta_4 + \frac{55552}{81} \zeta_5 + \frac{1792}{9} \zeta_3^2 - \frac{3200}{9} \zeta_6 \right]$$

$$+ n_f^2 C_A C_F^2 \left[\frac{332254}{2187} - \frac{85016}{243} \zeta_3 + \frac{20752}{27} \zeta_4 - \frac{28544}{81} \zeta_5 - \frac{13952}{27} \zeta_3^2 + \frac{1600}{27} \zeta_6 \right]$$

$$+ n_f^2 C_A^2 C_F \left[\frac{631400}{6561} + \frac{214268}{243} \zeta_3 - 784 \zeta_4 - \frac{53344}{243} \zeta_5 + \frac{25472}{81} \zeta_3^2 + \frac{22400}{81} \zeta_6 \right]$$

$$- n_f^2 \frac{d_{FF}^{(4)}}{N_F} \left[\frac{43744}{81} - \frac{35648}{81} \zeta_3 - \frac{1792}{9} \zeta_4 - \frac{52480}{81} \zeta_5 + \frac{2048}{27} \zeta_3^2 + \frac{12800}{27} \zeta_6 \right]$$

$$+ n_f^3 C_F^2 \left[\frac{265510}{19683} + \frac{11872}{729} \zeta_3 - \frac{128}{3} \zeta_4 + \frac{512}{27} \zeta_5 \right]$$

$$+ n_f^3 C_A C_F \left[\frac{168677}{19683} + \frac{11872}{729} \zeta_3 + \frac{2752}{81} \zeta_4 - \frac{4096}{81} \zeta_5 \right] - n_f^4 C_F \left[\frac{5504}{19683} + \frac{1024}{729} \zeta_3 - \frac{128}{81} \zeta_4 \right],$$

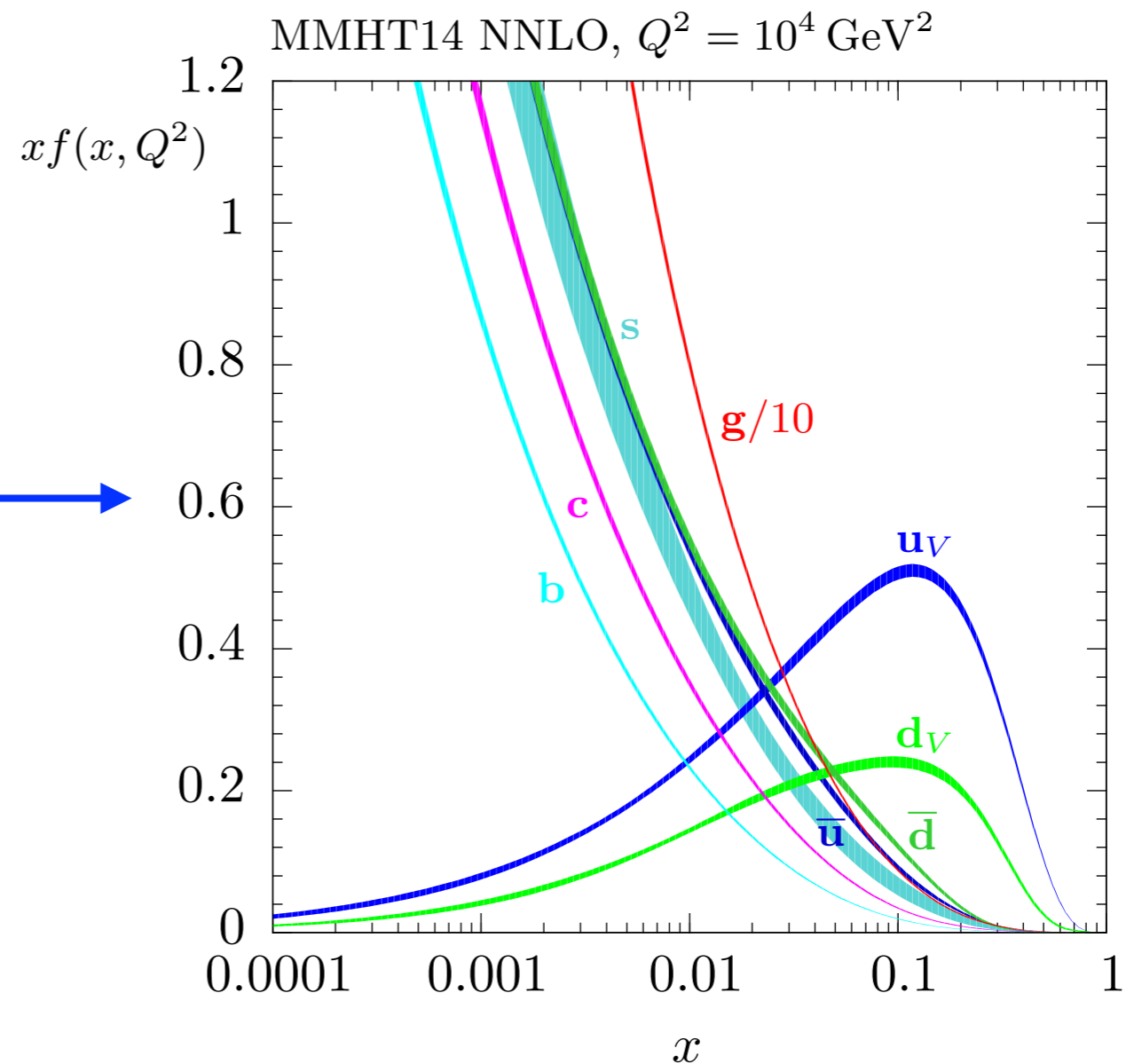
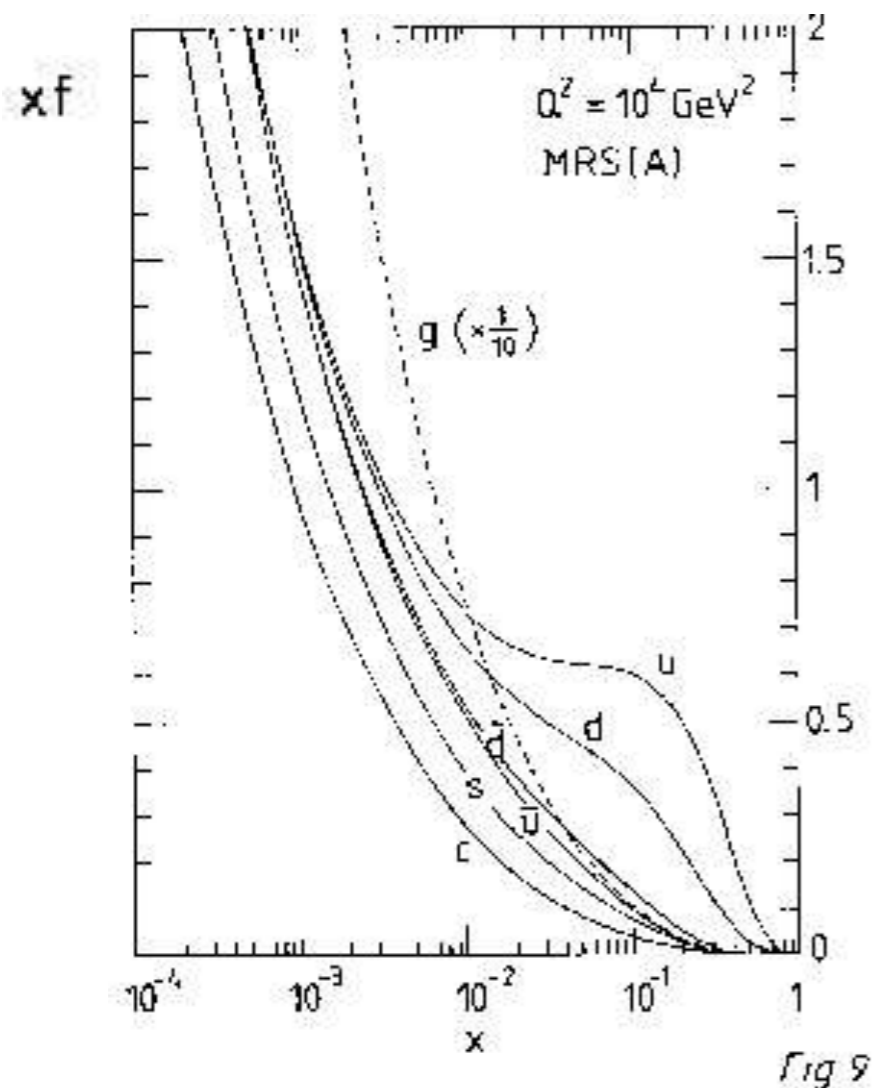
+ 1.5 more pages!

Where are we now?

- Data and theory at **unprecedented level of precision** provide unique opportunity to really pin down proton structure like never before.

~ 2 decades after quark discovery.
Still only rough picture.

Today: Precision Physics!



- Taking a closer look...

The Proton Backbone

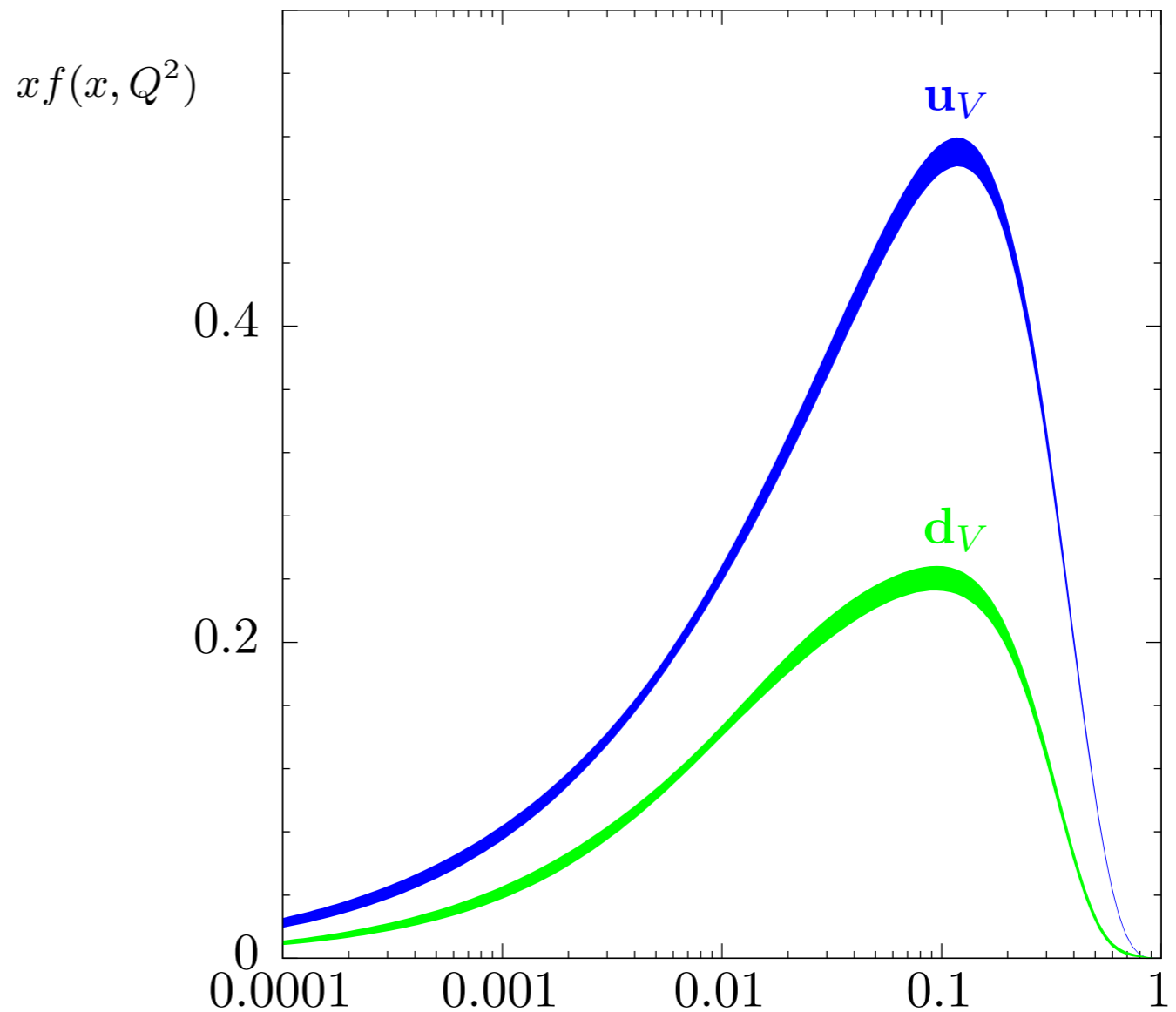
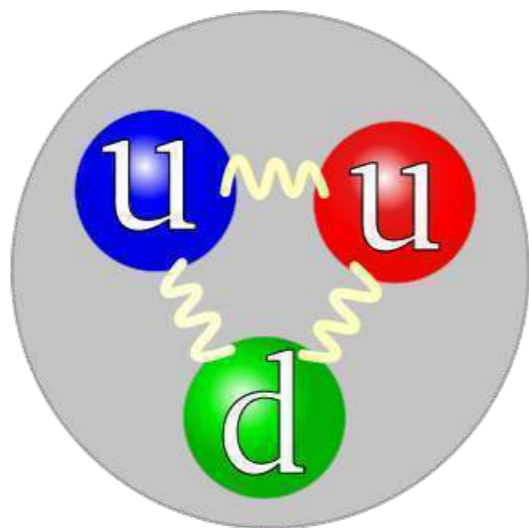
- Within our precise mapping, many features present...

$$u - \bar{u}, d - \bar{d}$$

- ‘Valence’ up and down quark structure consistent with basic **uud** picture.

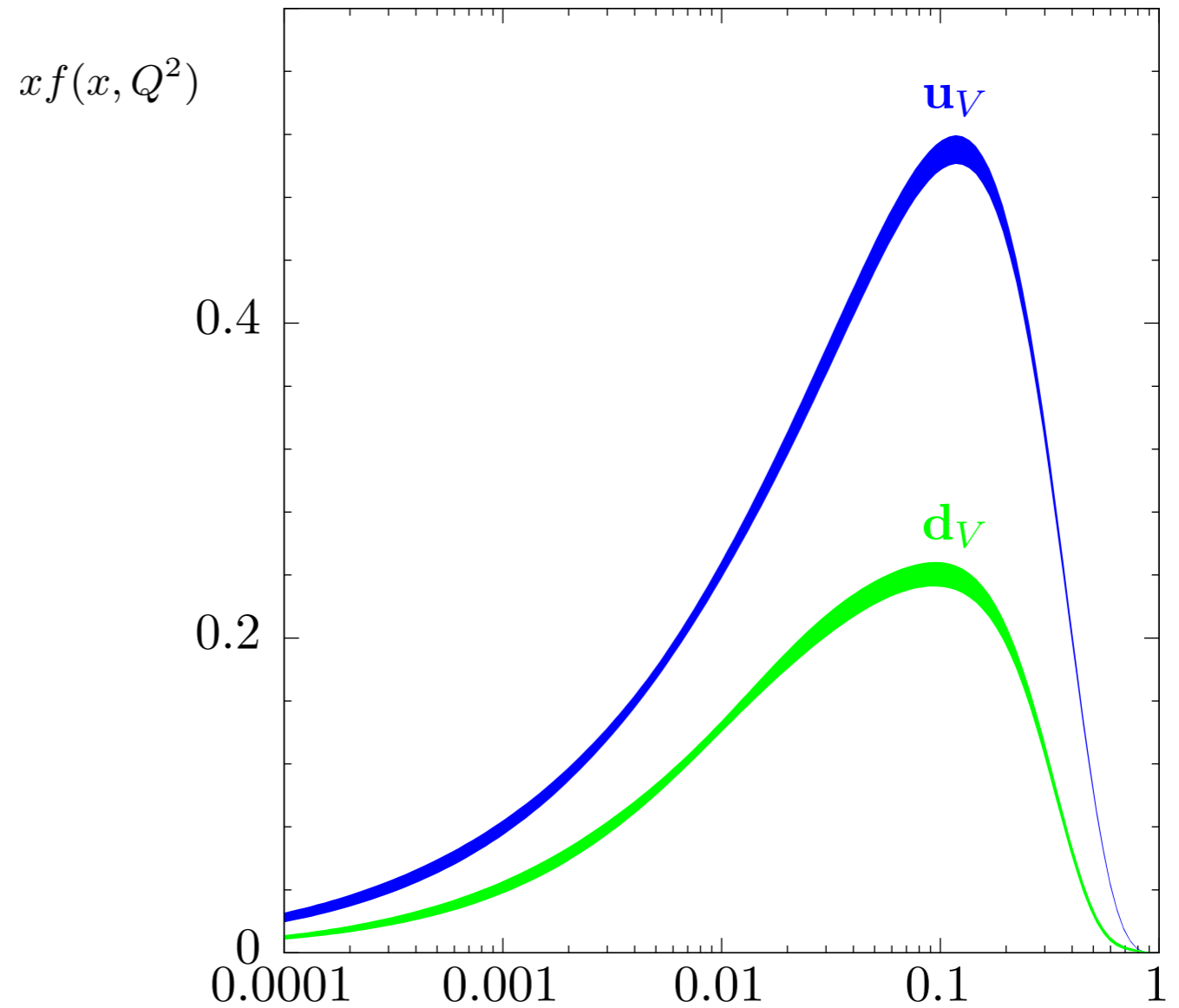
- $u_V \approx 2d_V$

- Peaked at $x \approx \frac{1}{3}$.



Glucos

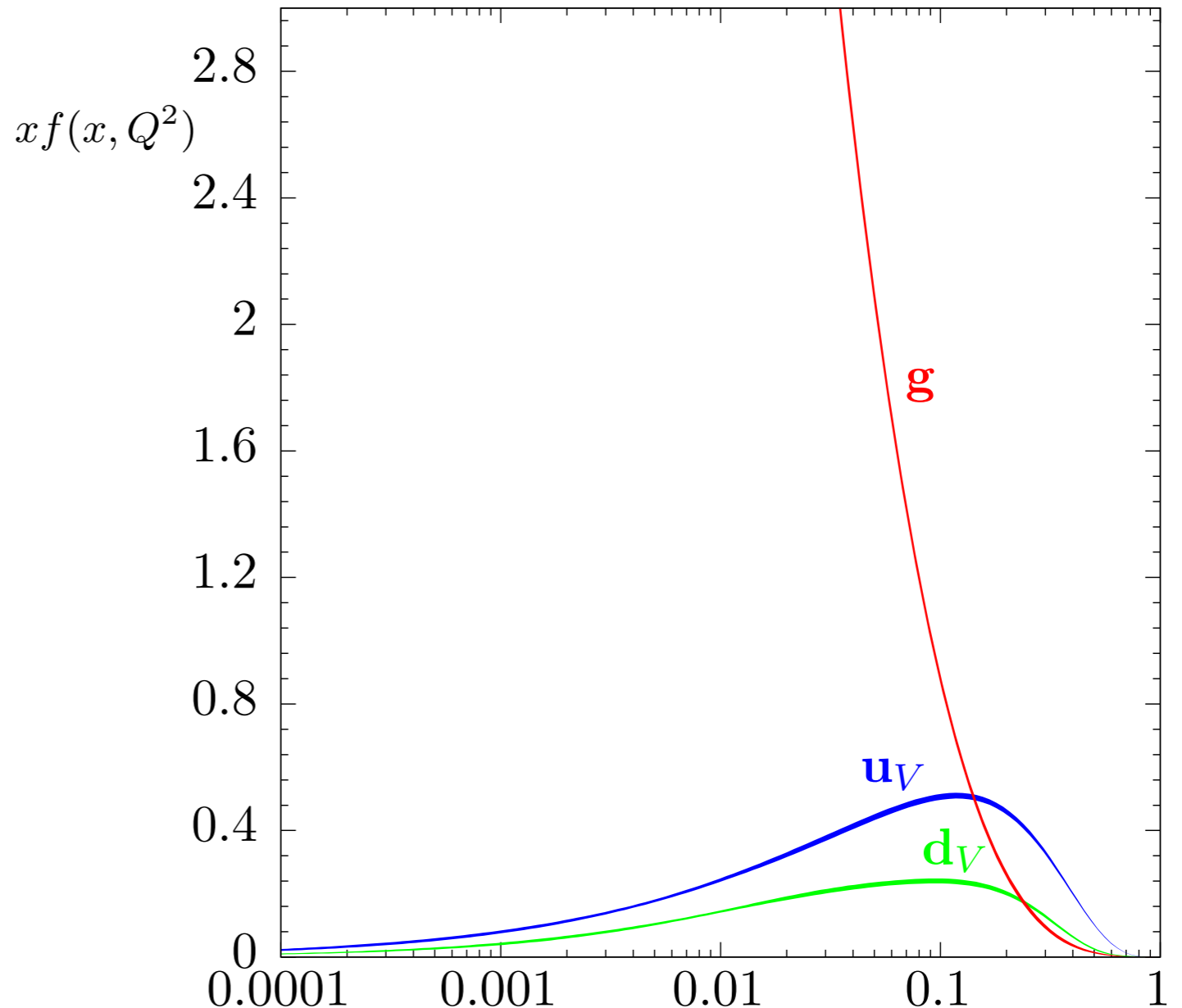
- What happens when we add gluons into the picture?



Gluons!

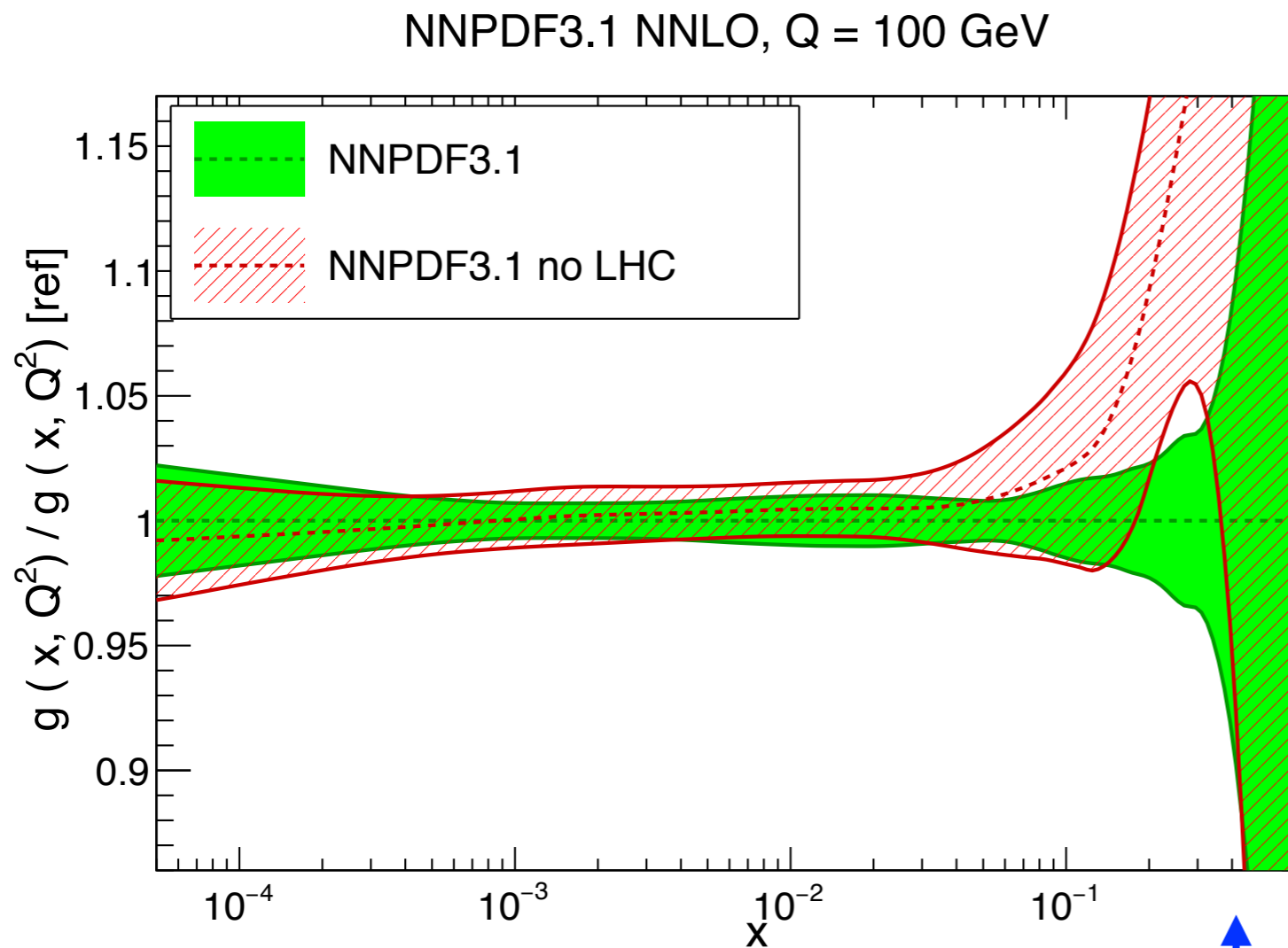
- Contribution from gluons is absolute huge!
- In fact, roughly **50%** of proton energy carried by gluons.
- LHC a gluon-gluon collider: crucial for Higgs physics.

(Fabrizio 's talk)



The Gluon: Zooming In

- **Roughly** 50% of proton energy carried by gluons. Research today: pinning this statement down with **high precision**.



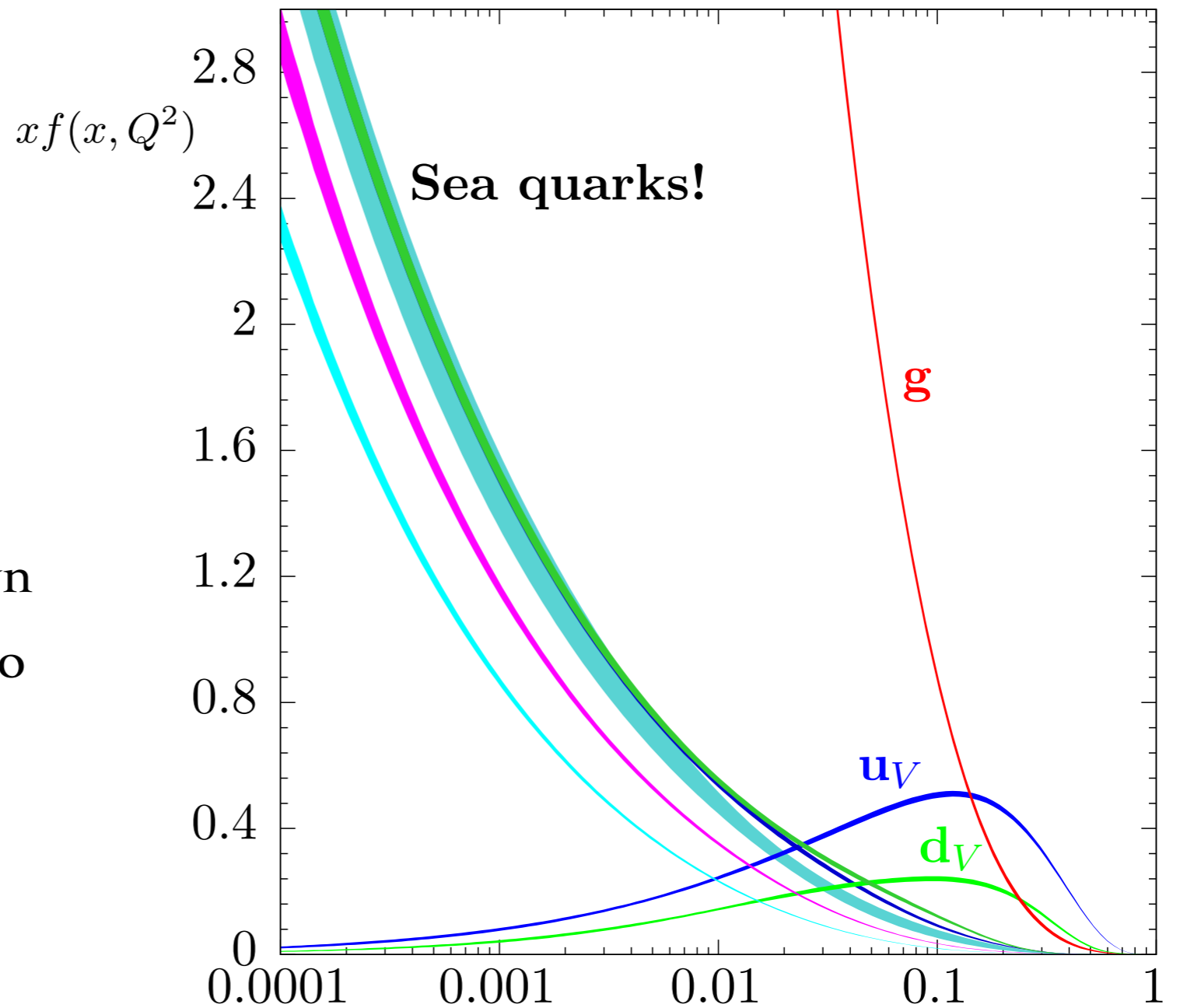
- Modern fits now achieve **~ 1%** level precision in many regions.
- Not always the case. Key question: how can we best exploit **LHC** to push into less understood regions?

Higgs

Physics beyond the
Standard Model

The Quark Sea

- Quark sea completely dominant over 'valence' uud in many regions.
- Crucial to pin these down precisely- direct relevant to LHC physics.



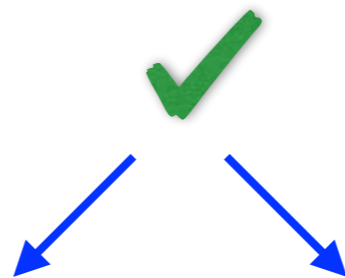
What next?

- Data and theory at **unprecedented level of precision** provide unique opportunity to really pin down proton structure. But also a challenge:
 - ❖ Dealing with **complex**/computationally expensive **theory** effectively.
 - ❖ Dealing with cases where data and theory **do not agree** to the high precision standard we require.
 - ❖ Re-evaluating foundations: are PDF uncertainties **accurate** as well as precise? Other ‘theory’ **uncertainties/biases** hidden in fits?
 - ❖ Perhaps we can calculate PDFs after all!? Applying numerical ‘**lattice**’ techniques to predict PDFs, rather than fitting them.
- ...and much more besides. Ensuring that our knowledge of PDFs matches requirements for understanding in detail the Higgs at the LHC.

Higgs Production at the LHC

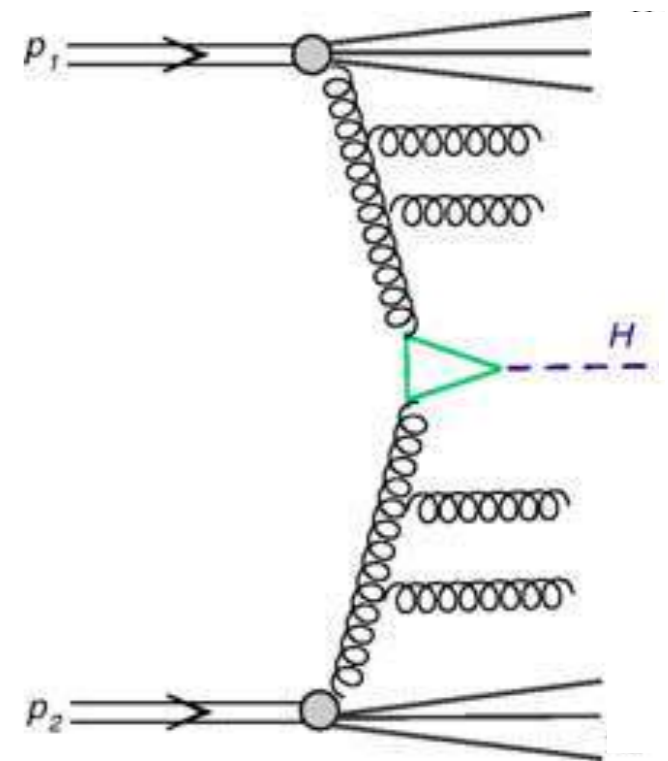
- We have all we need to go from protons to collisions at the **LHC**.
- Key idea: LHC is not really a proton-proton collider. It is a quark-quark, gluon-gluon... collider.

Fabrizio



$$\sigma(pp \rightarrow h + X) \sim \sigma(gg \rightarrow h) \otimes g(x_1, Q^2) \otimes g(x_2, Q^2) ,$$

- **Next question:** how do we predict the rate of Higgs boson production from quark-quark, gluon-gluon... collisions? On to Fabrizio...



Thank you for listening!